

Integrated waveguide photonics circuits for quantum simulation and beyond

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<http://quantumoptics.phys.uniroma1.it>

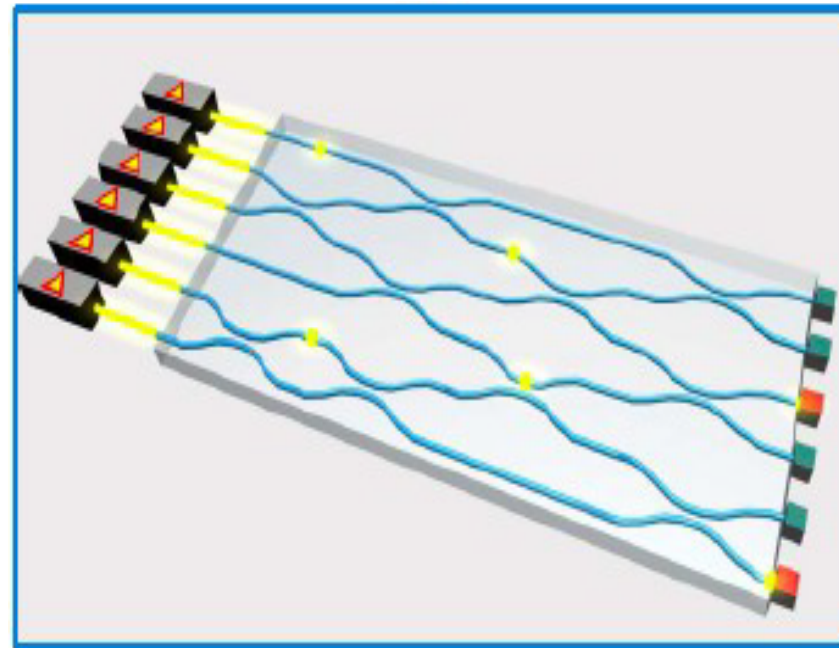
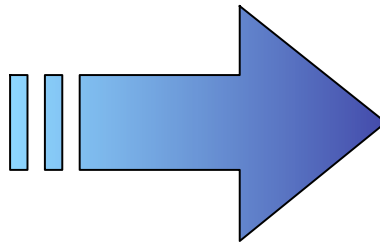
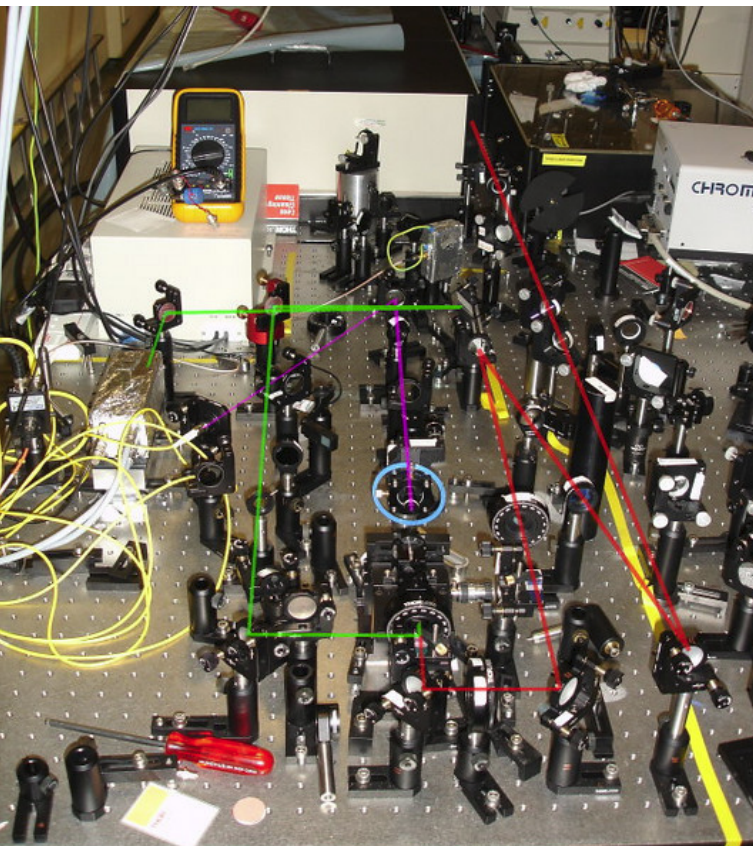
Dipartimento di Fisica Pavia, 25 Febbraio 2016

Integrated photonic quantum circuits

A new scenario for photonic quantum information processing

Complex optical setup affected by:

- ✓ Large physical size
- ✓ Low stability
- ✓ Difficulty to move forward applications outside laboratory



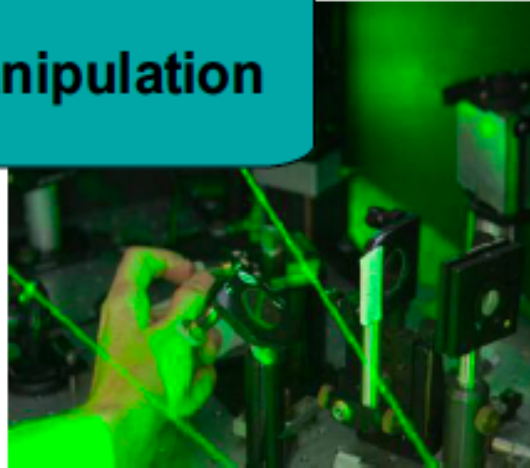
A. Politi et al., Science 320, 646 (2008)

On-chip integrated setup

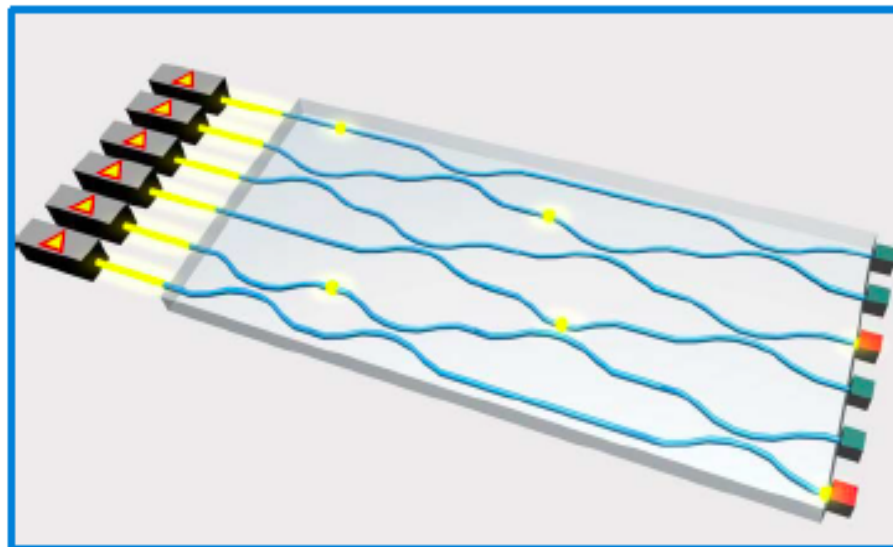
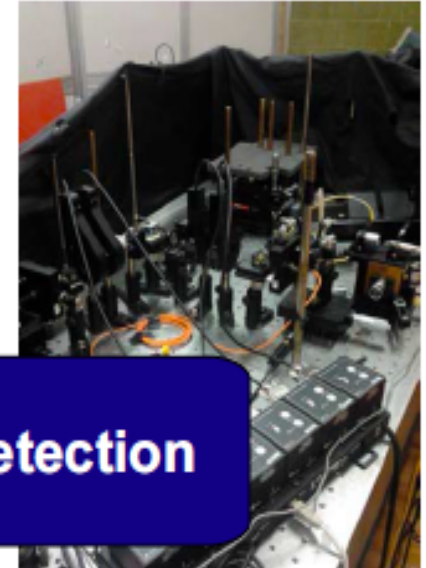
Preparation



Manipulation

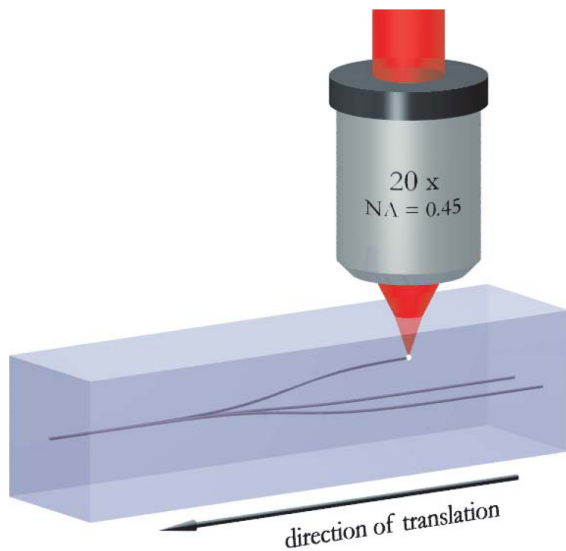


Detection



Femtosecond laser writing

3-D capability

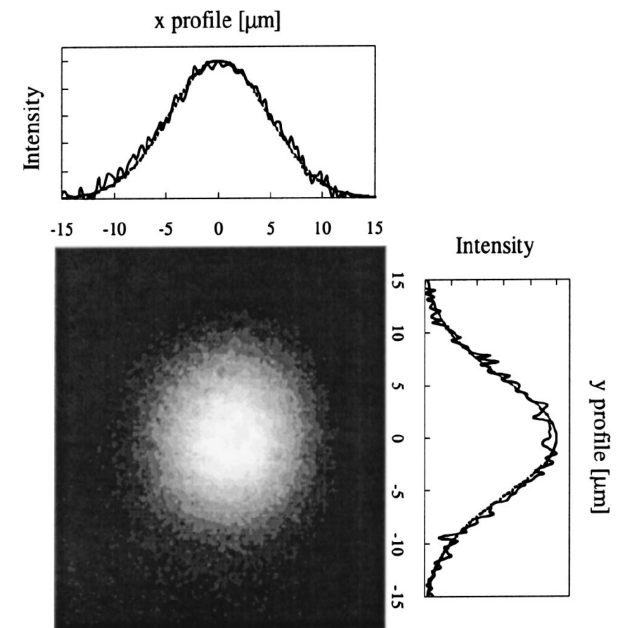


Rapid device prototyping:
writing speed = 4 cm/s

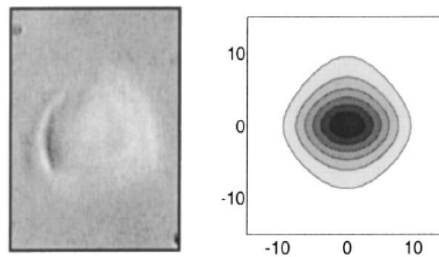
Propagation of circular
gaussian modes

Main features

Circular waveguide
transverse profile



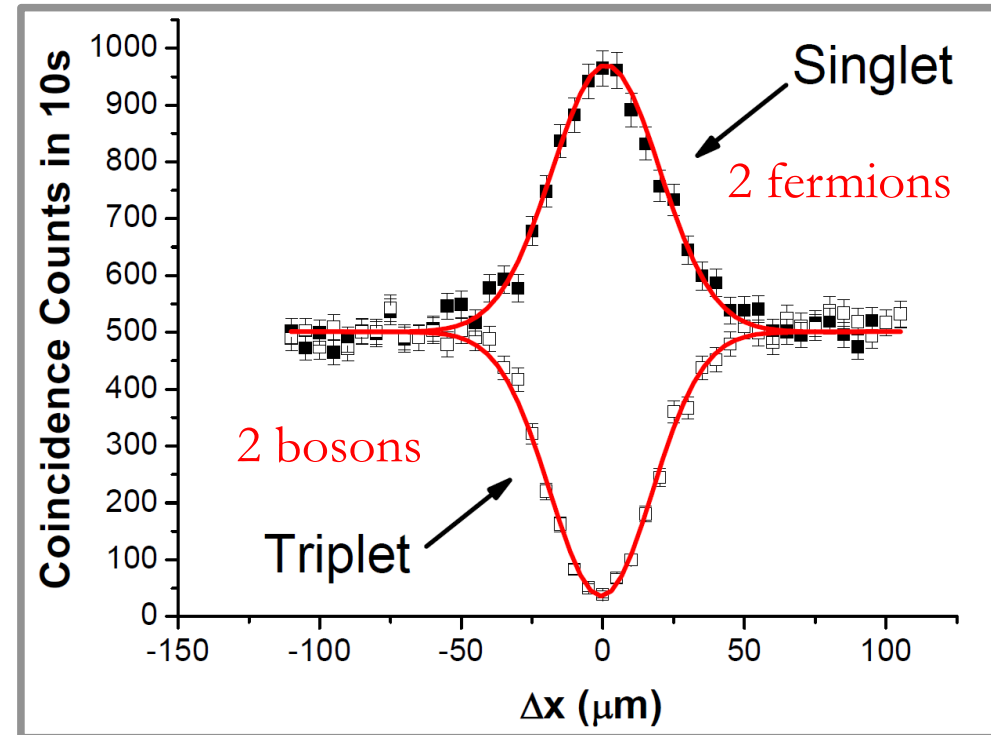
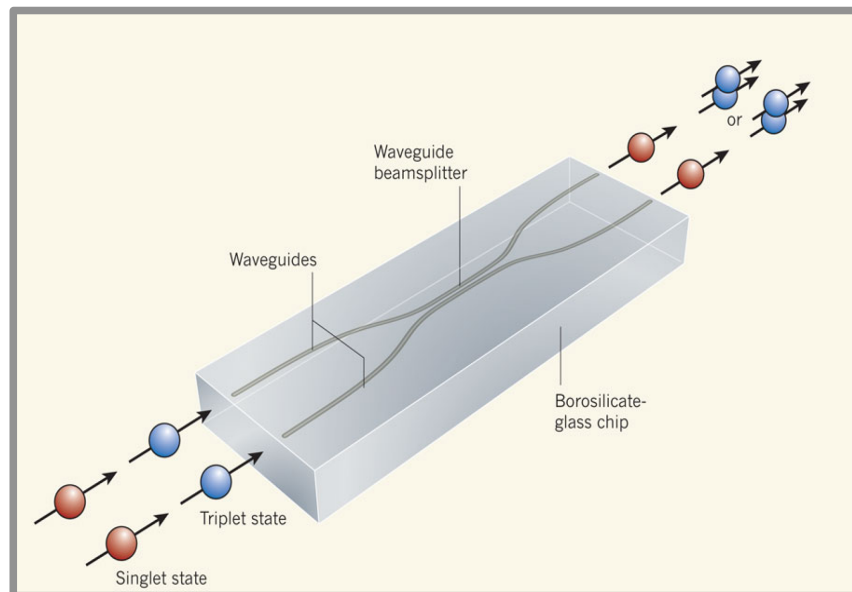
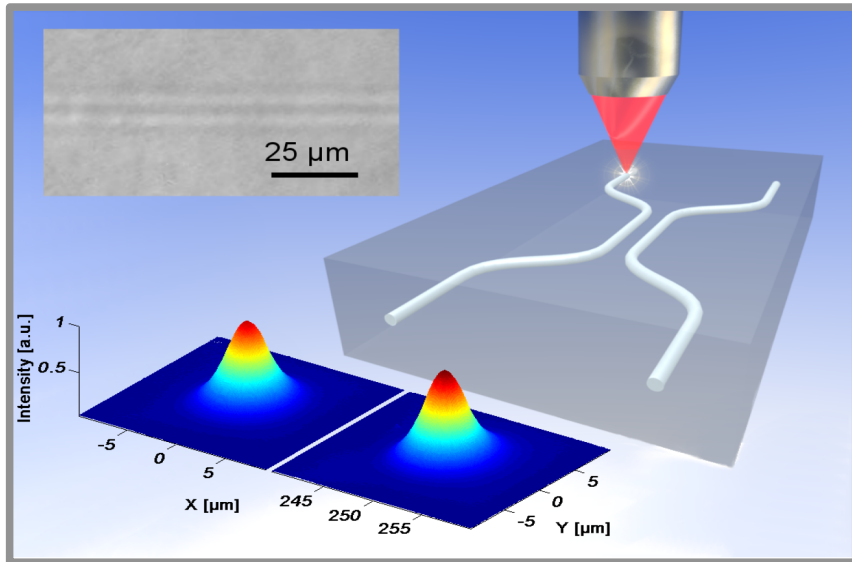
Low birefringence



ABLE TO SUPPORT ANY POLARIZATION STATE

On-chip integrated directional coupler

$$|\Psi^\phi\rangle = \frac{1}{\sqrt{2}}(|H\rangle_A|V\rangle_B + e^{i\phi}|V\rangle_A|H\rangle_B)$$



L. Sansoni *et al.* *Phys. Rev. Lett.* **105**, 200503 (2010)

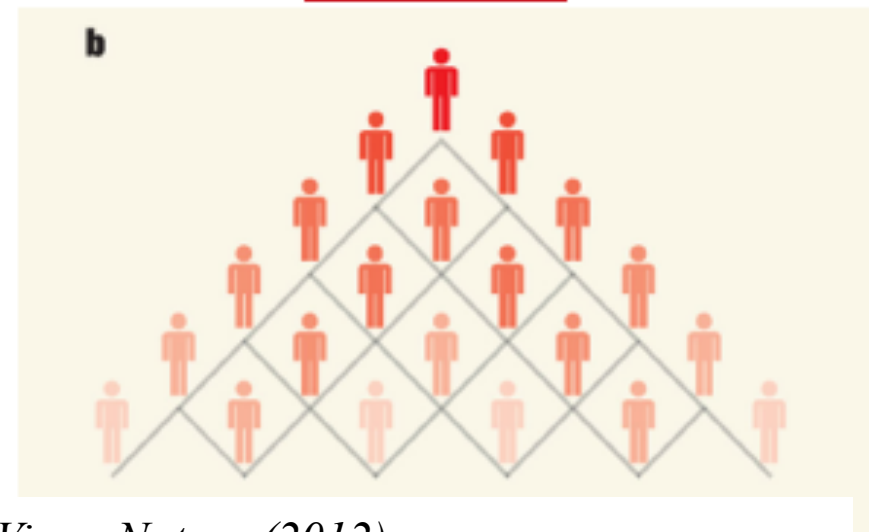
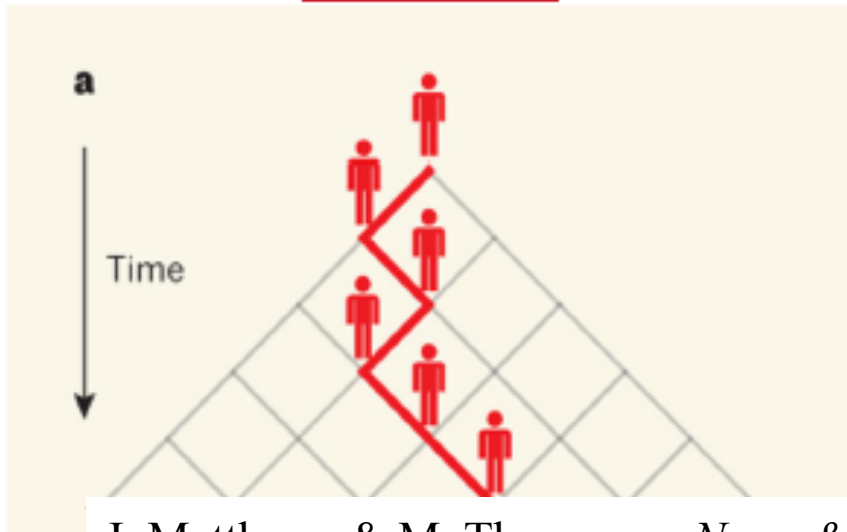
Goal: simulate the dynamics of non-interacting bosons/fermions in complex interferometric structures

Discrete time Quantum Walk

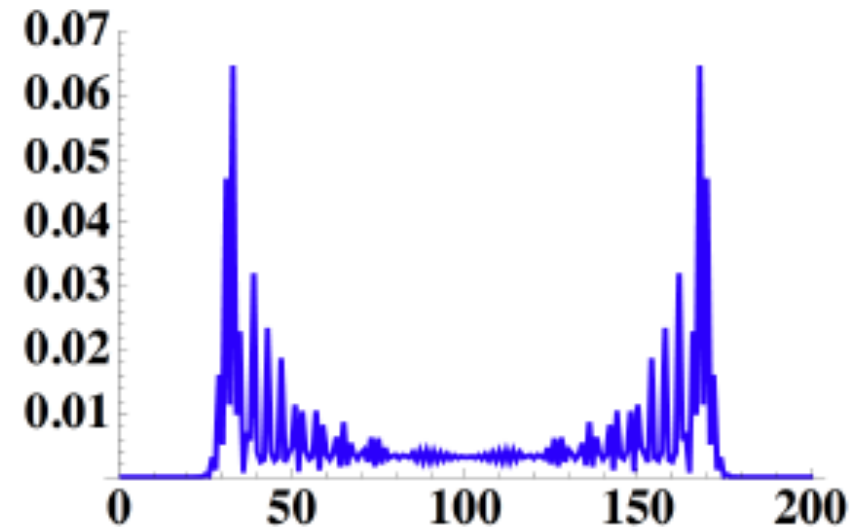
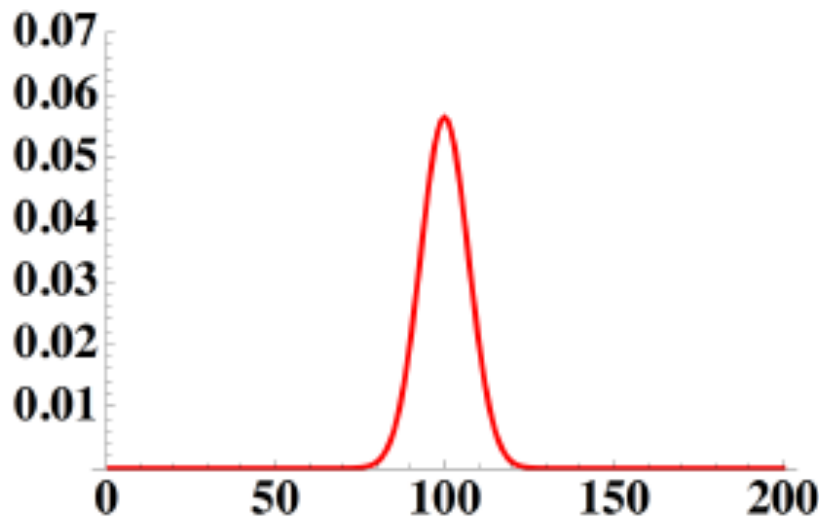
Quantum walk: extension of the classical random walk:
a walker on a lattice “jumping” between different sites with given probability

Classical

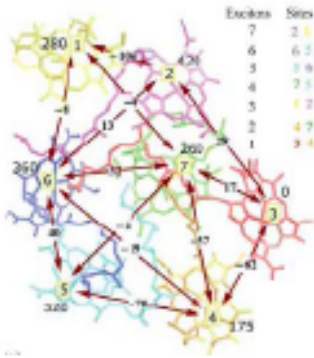
Quantum



J. Matthews & M. Thompson, *News & Views Nature* (2012)

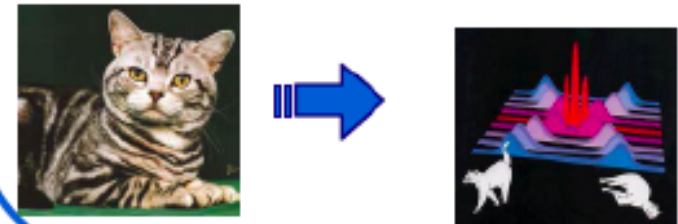


Why Quantum Walks?

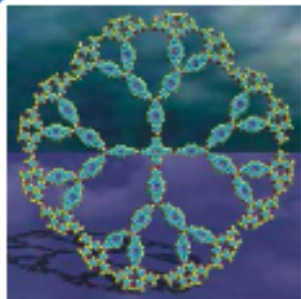


Energy transfer:
within photosynthetic systems can display quantum effects such as delocalized excitonic transport which can be simulated by QW.

Controlled transition from Classical to Quantum:
QW can be employed for testing the transition from the quantum to the classical world by applying a controlled degree of decoherence.

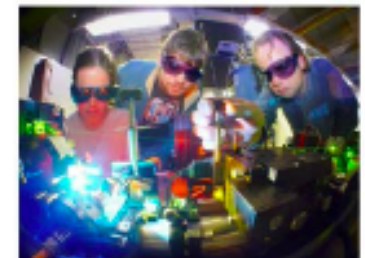


QW?



Light-harvesting molecule
is efficient at concentrating light at its center as quantum walk reaches the target vertex exponentially faster than a classical walk: because of destructive interference between the paths that point backward, toward the leaves.

Faster quantum Computation:
It has been theoretically proven that QWs allow the speed-up of search algorithms



Discrete time Quantum Walk



In a discrete Quantum Walk one or more quantum particles evolve on a graph with their evolution governed by their internal quantum coin (QC) states

Walker in position j described by the quantum coin state $|j\rangle$. The particle shifts up or down depending on the internal QC state $|U\rangle$ or $|D\rangle$

Walk evolution described by the operator:

$$E = \sum_j |j-1\rangle\langle j| \otimes |U\rangle\langle U| + |j+1\rangle\langle j| \otimes |D\rangle\langle D|$$

Experimental platforms:

Ion trap

F. Zahringer *et al.*,
PRL **104** 100503 (2010)

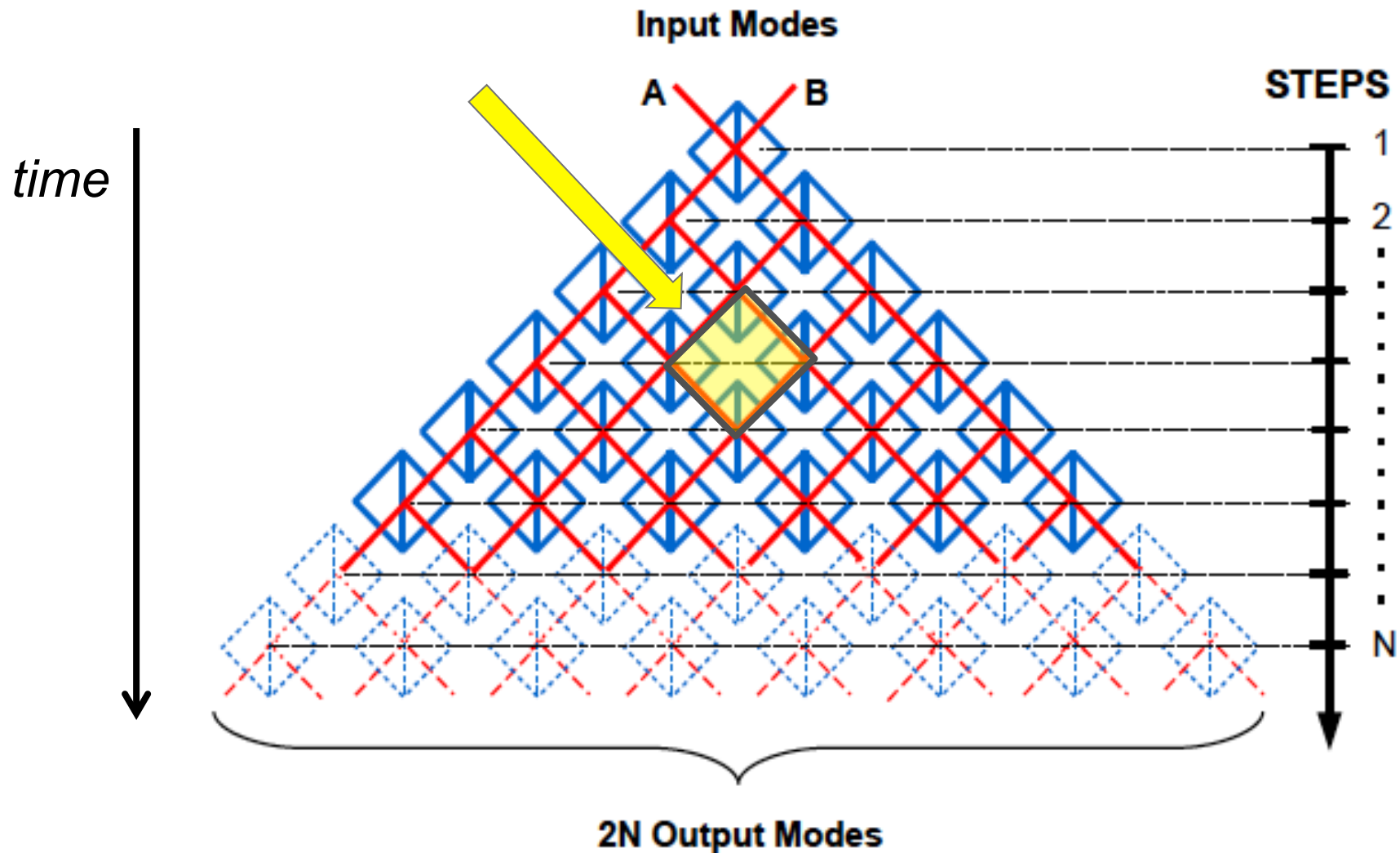
Fiber loops

A. Schreiber *et al.*,
PRL **104** 050502 (2010)
2D QW: Science (2012)

Coupled waveguides

A. Peruzzo *et al.*,
Science **329** 1500 (2010)

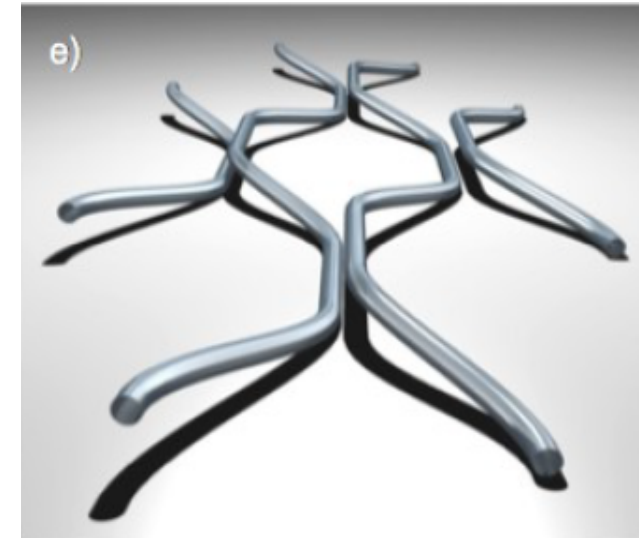
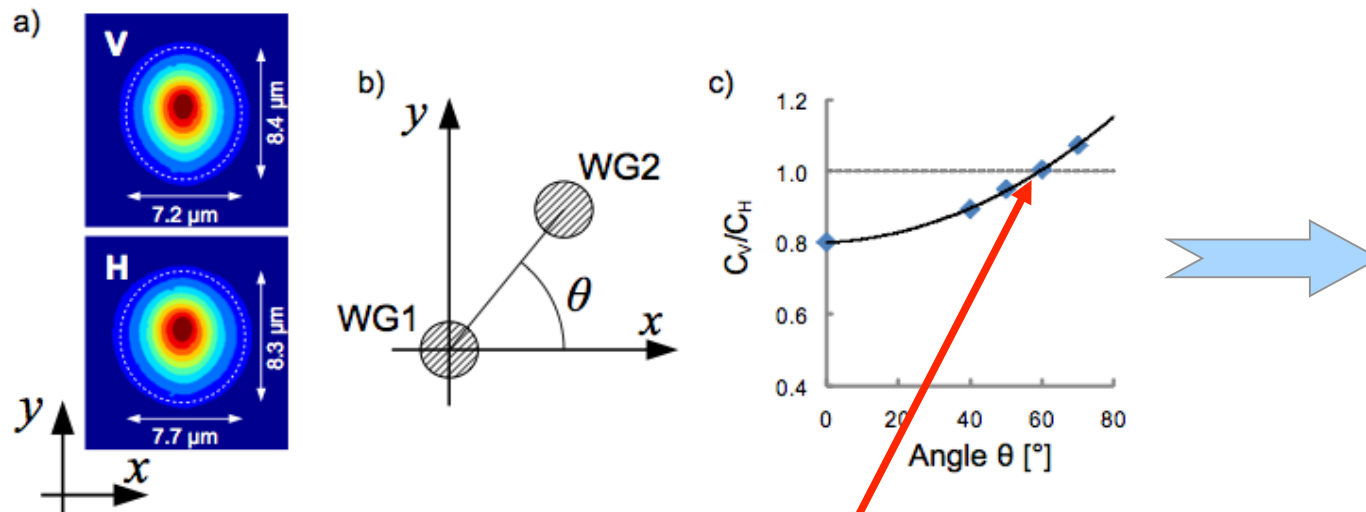
Photonic QW



Feasible with integrated photonic waveguides

Polarization independent QW

V-mode slightly larger than the H-mode along the y axis:
Use 3D capability to improve the on-chip devices



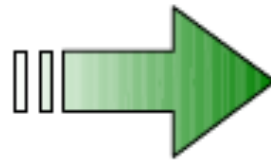
$$\mathbf{R}_H = \mathbf{R}_V$$

- 3D array of beamsplitters with balanced reflectivities $\mathbf{R}_H = \mathbf{R}_V = 49\%$ able to support any polarization state.
- Path lengths controlled up to few nanometers: complete control of phase difference.

2-particle QW

The symmetry of two travelling quantum walkers influences the output probability distribution

Polarization independent integrated beam splitter



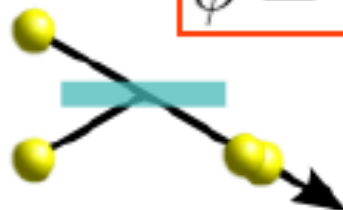
Exploit polarization entanglement to change the symmetry of two-particle wavefunction

$$|\Psi^\phi\rangle = \frac{1}{\sqrt{2}} (|H\rangle_A |V\rangle_B + e^{i\phi} |V\rangle_A |H\rangle_B)$$

Bosons

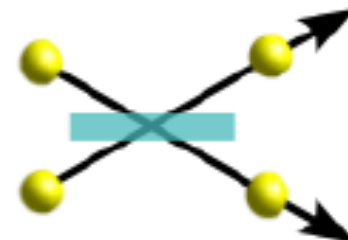


or



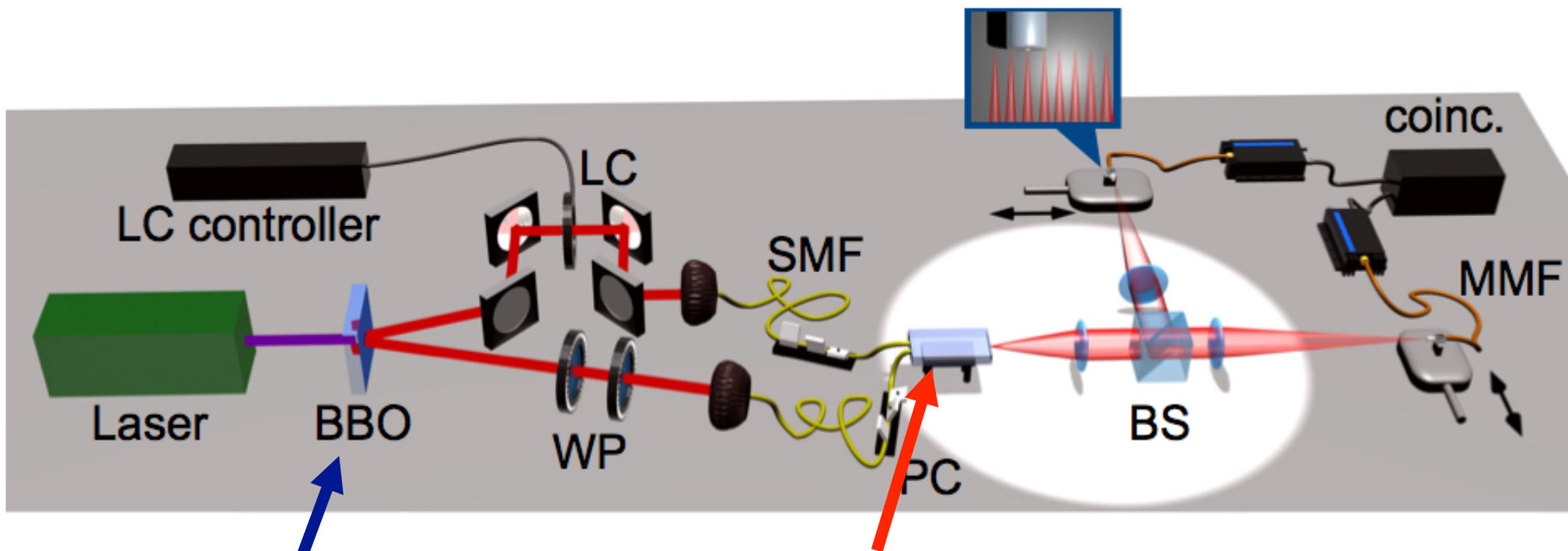
$$\phi = 0$$

Fermions

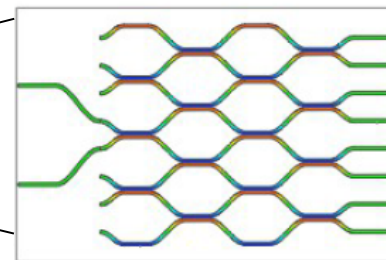
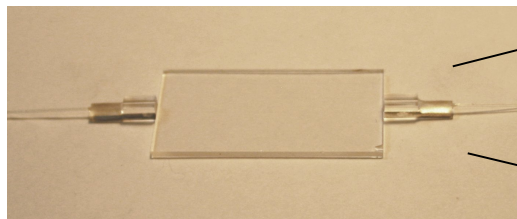


$$\phi = \pi$$

The experiment



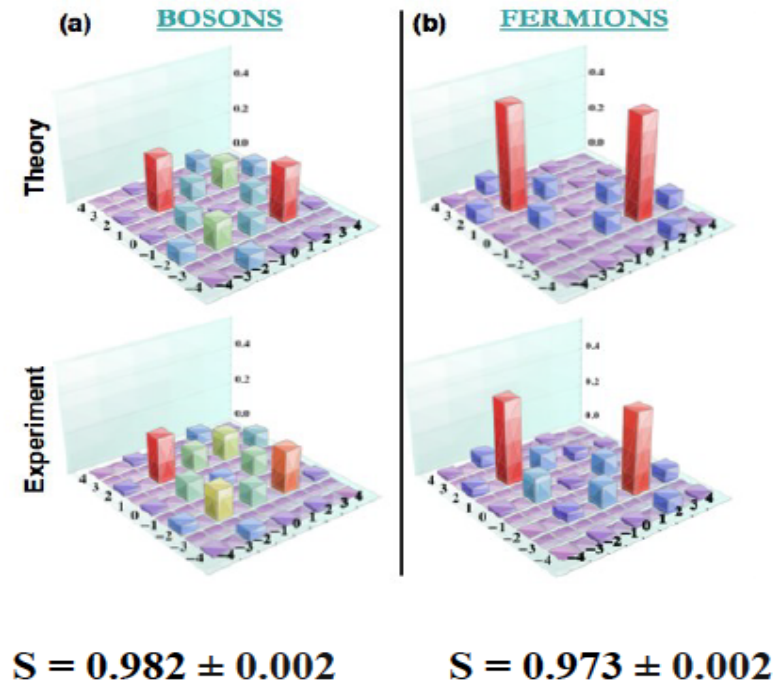
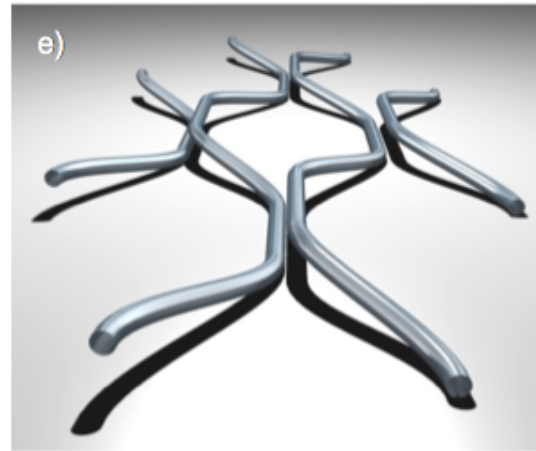
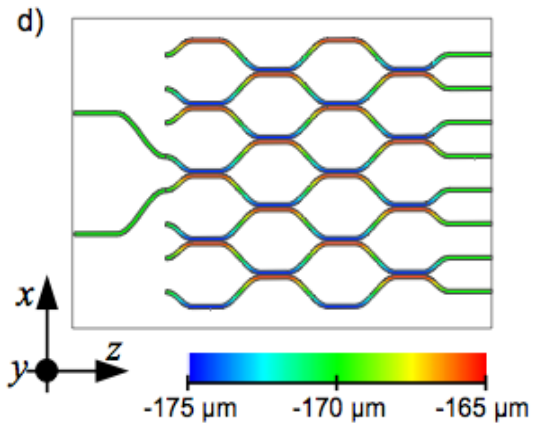
On chip BS array



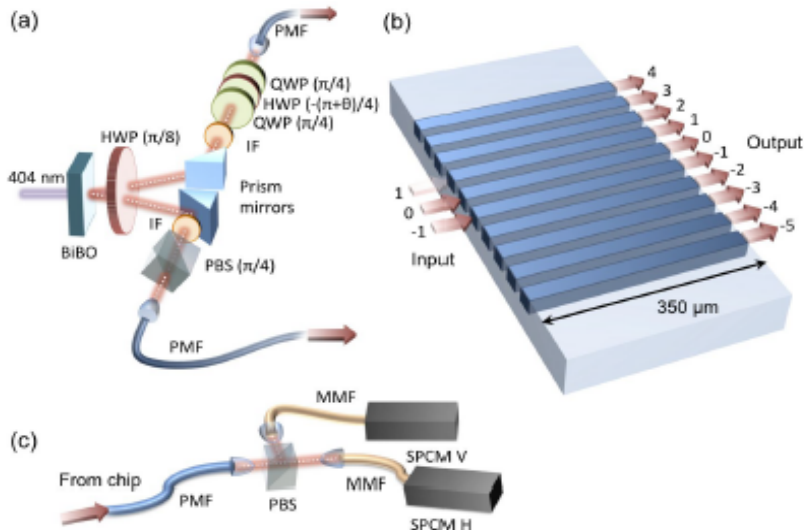
2-photon entangled states with different symmetries

$$|\Psi^\phi\rangle = \frac{1}{\sqrt{2}}(|H\rangle_A|V\rangle_B + e^{i\phi}|V\rangle_A|H\rangle_B)$$

Integrated Quantum Walk

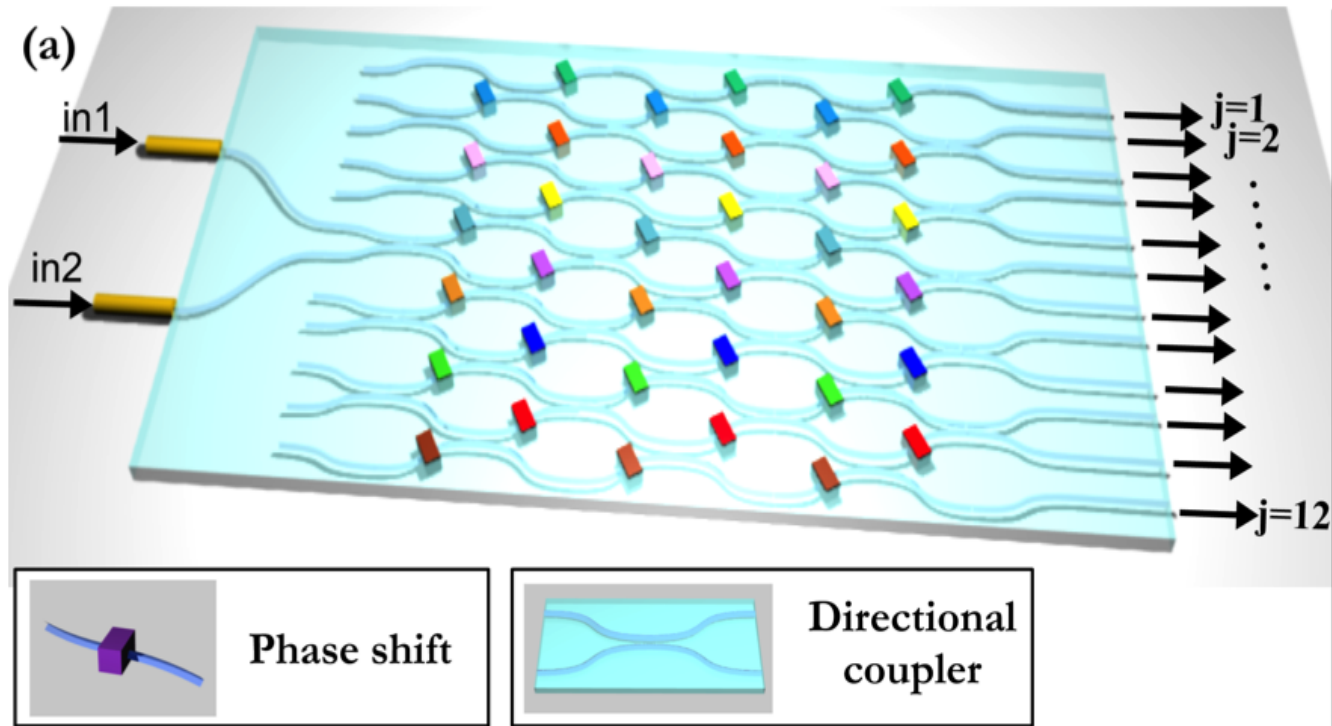


L. Sansoni et al. PRL (2012)
 Discrete time QW. Mimic bosonic/fermionic behaviour
 using two entangled photons

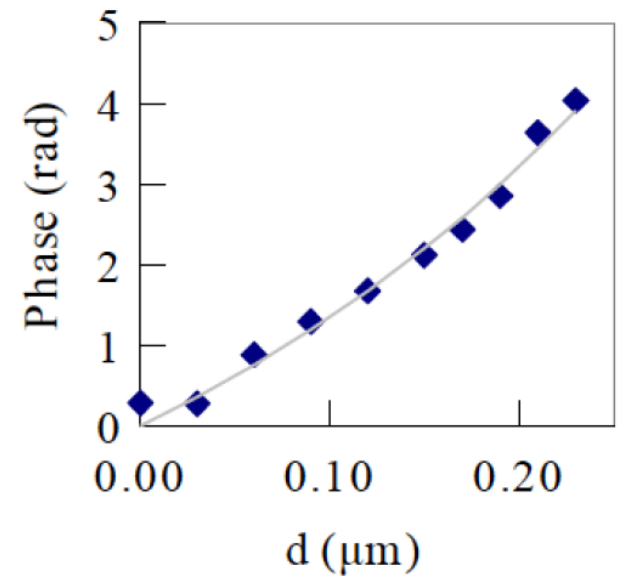
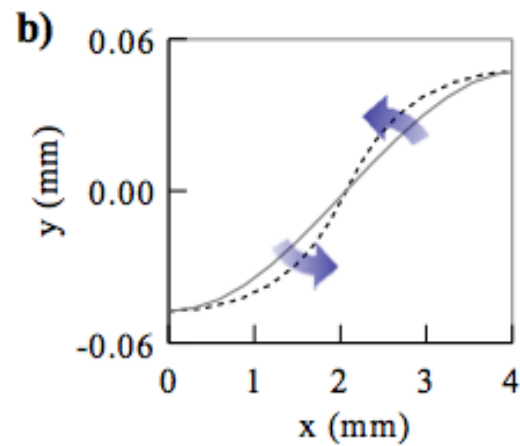
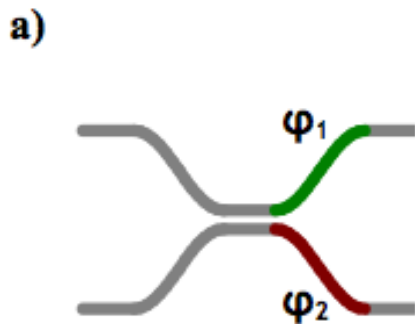


J.C.F. Matthews et al. Sci. Rep. (2013)
 Simulation of fermionic statistics

Disorder control in a QW



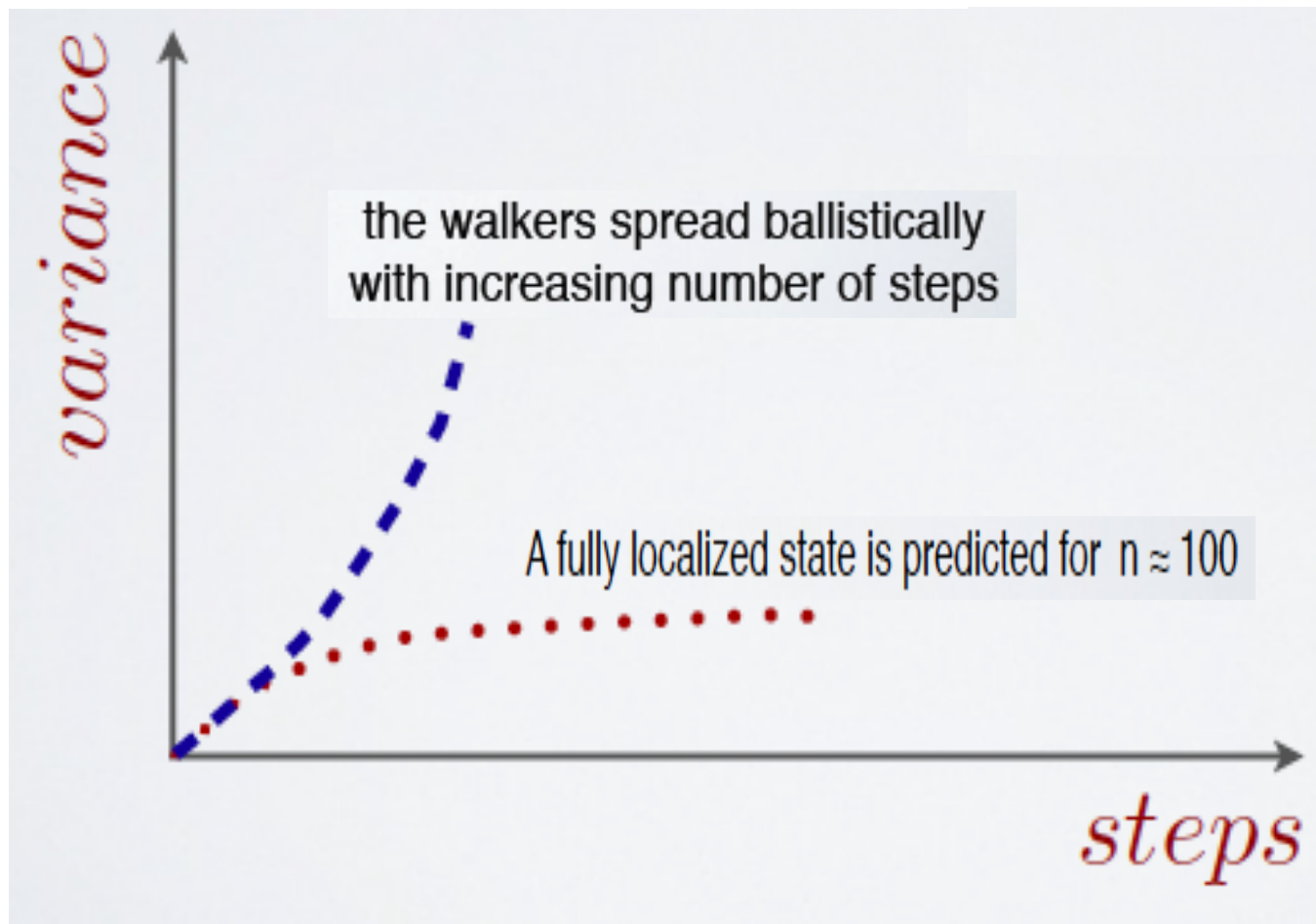
Phase shift induced by path variation



Ordered vs. disordered lattices

By realizing different phase maps **the pure role of symmetries** in different diffusion processes of two non-interacting bosons/fermions **may be investigated**.

Measure the variance of the distribution



Static disorder

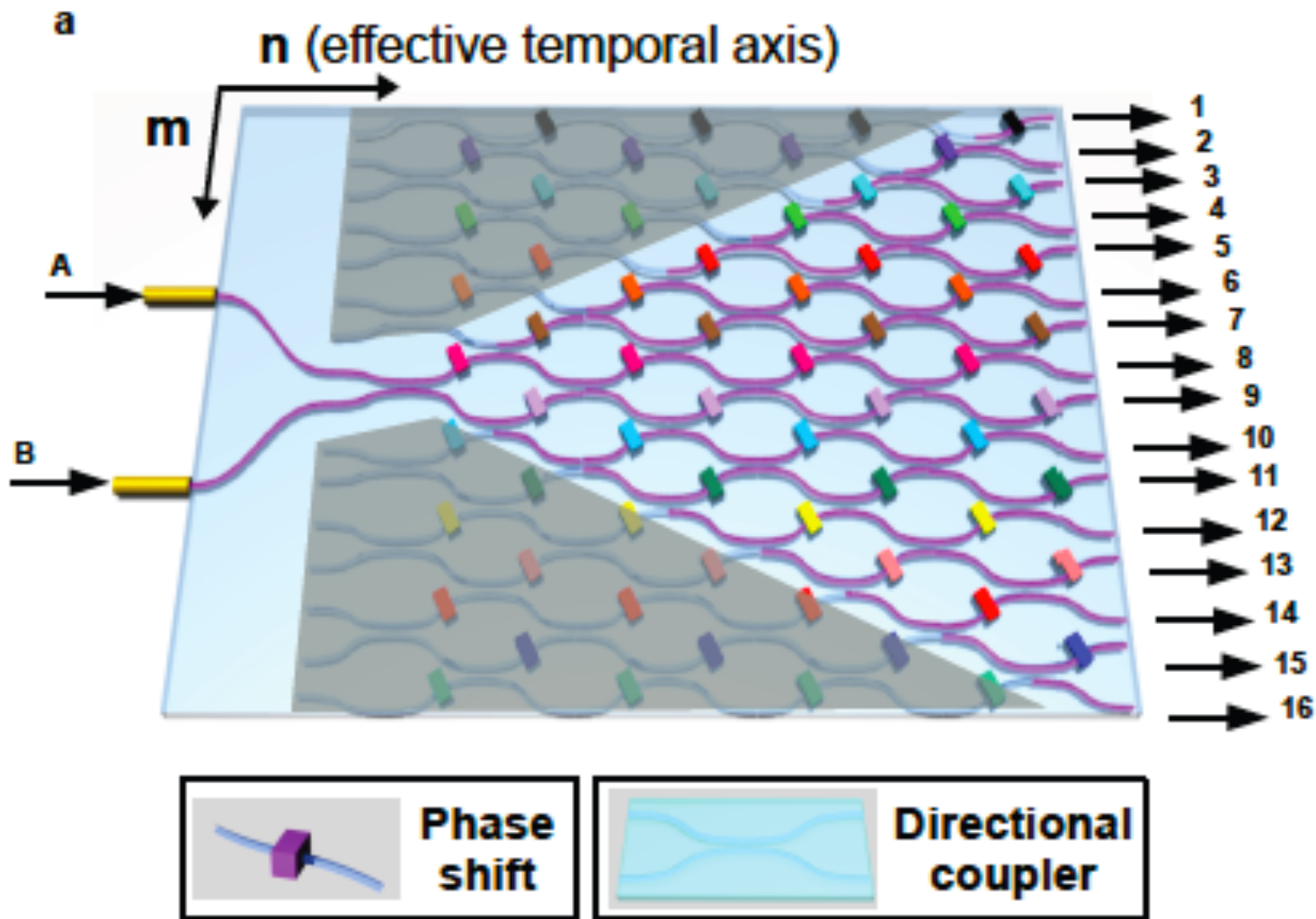
Disorder depending

- on site
- but **NOT** on time

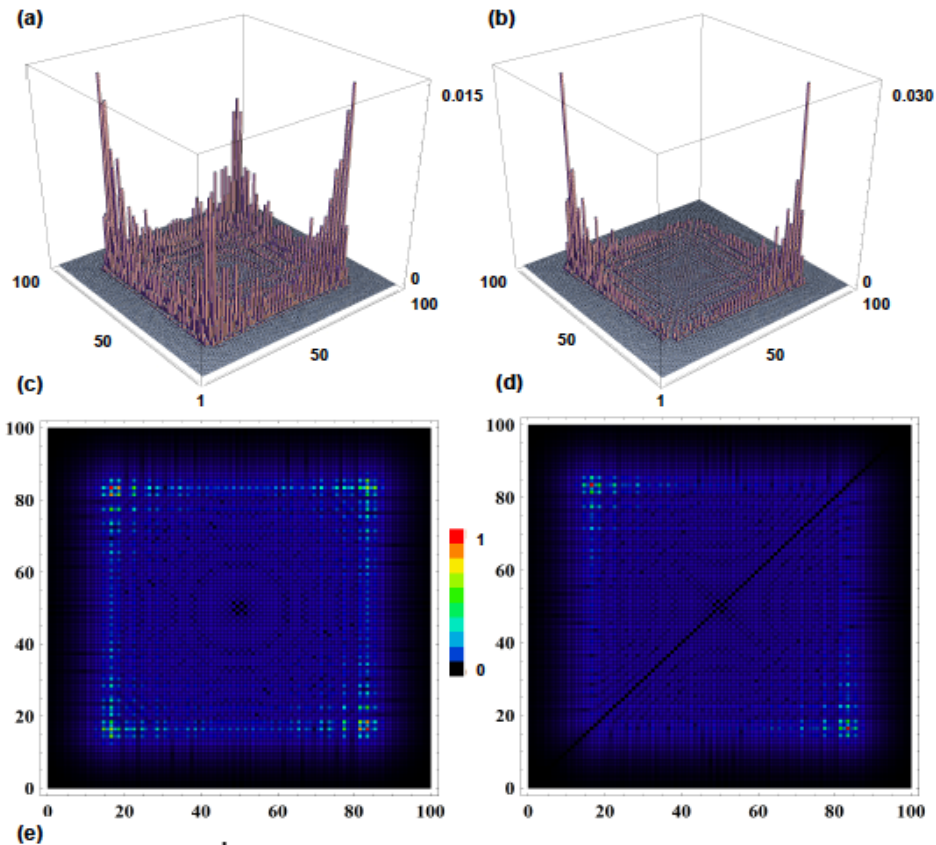


Anderson localization of the quantum particle wavefunction

Up to 64 polarization independent BSs and phase-shifters

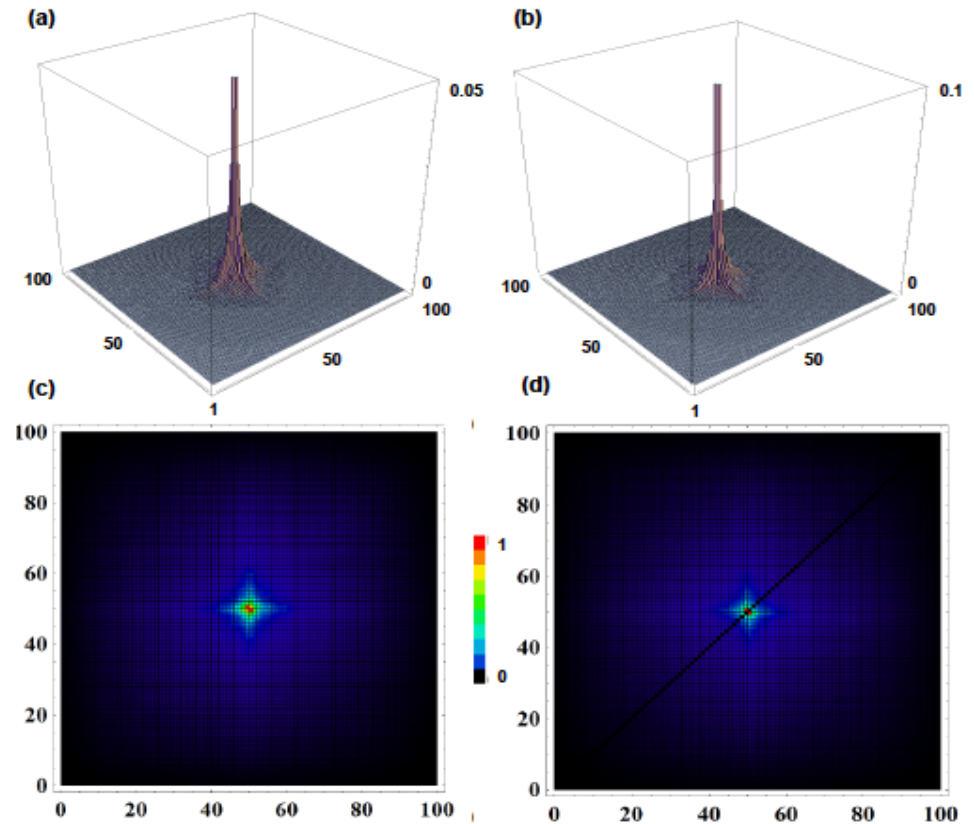


2-particle QW: theoretical distributions



2 Bosons
Ordered QW (50 steps)

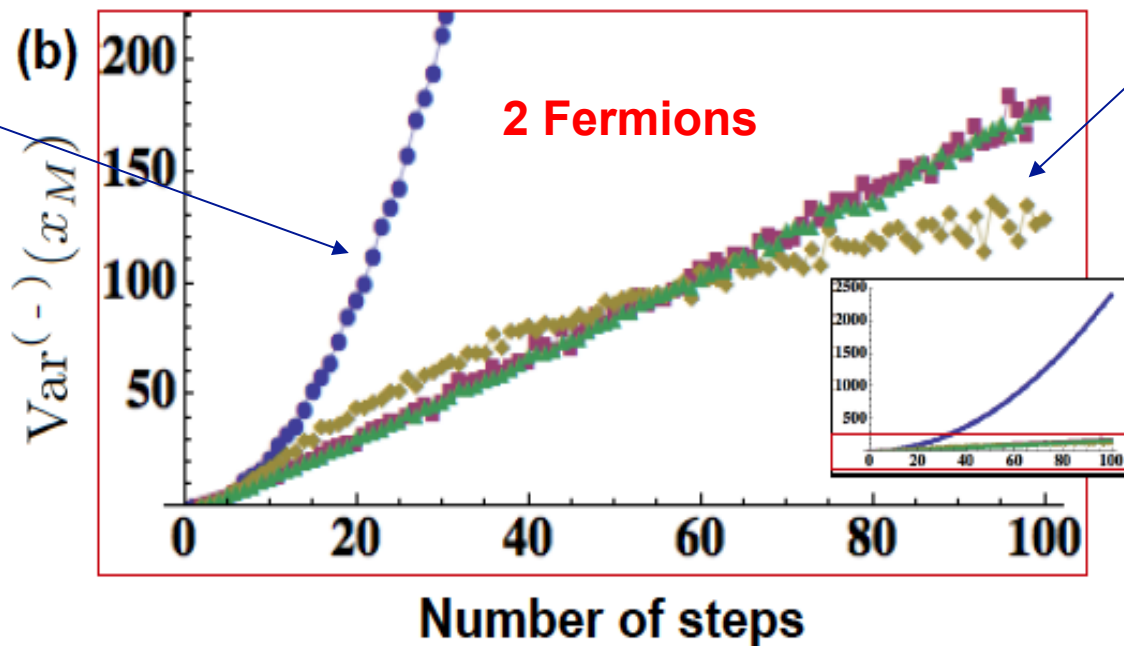
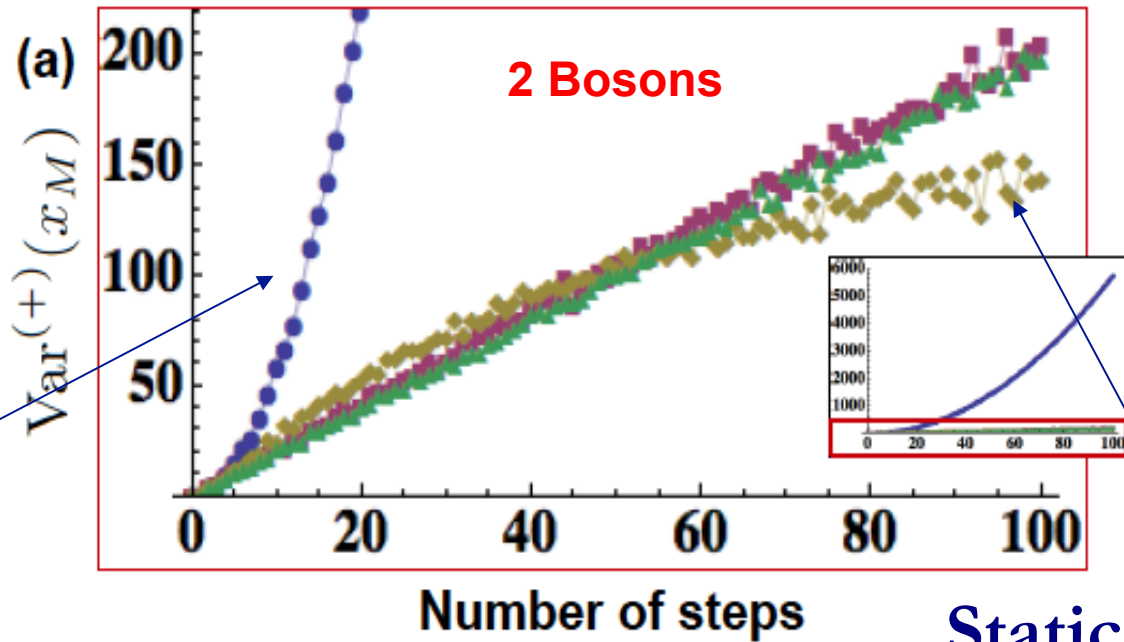
2 Fermions



2 Bosons
Static disordered QW (50 steps)

2 Fermions

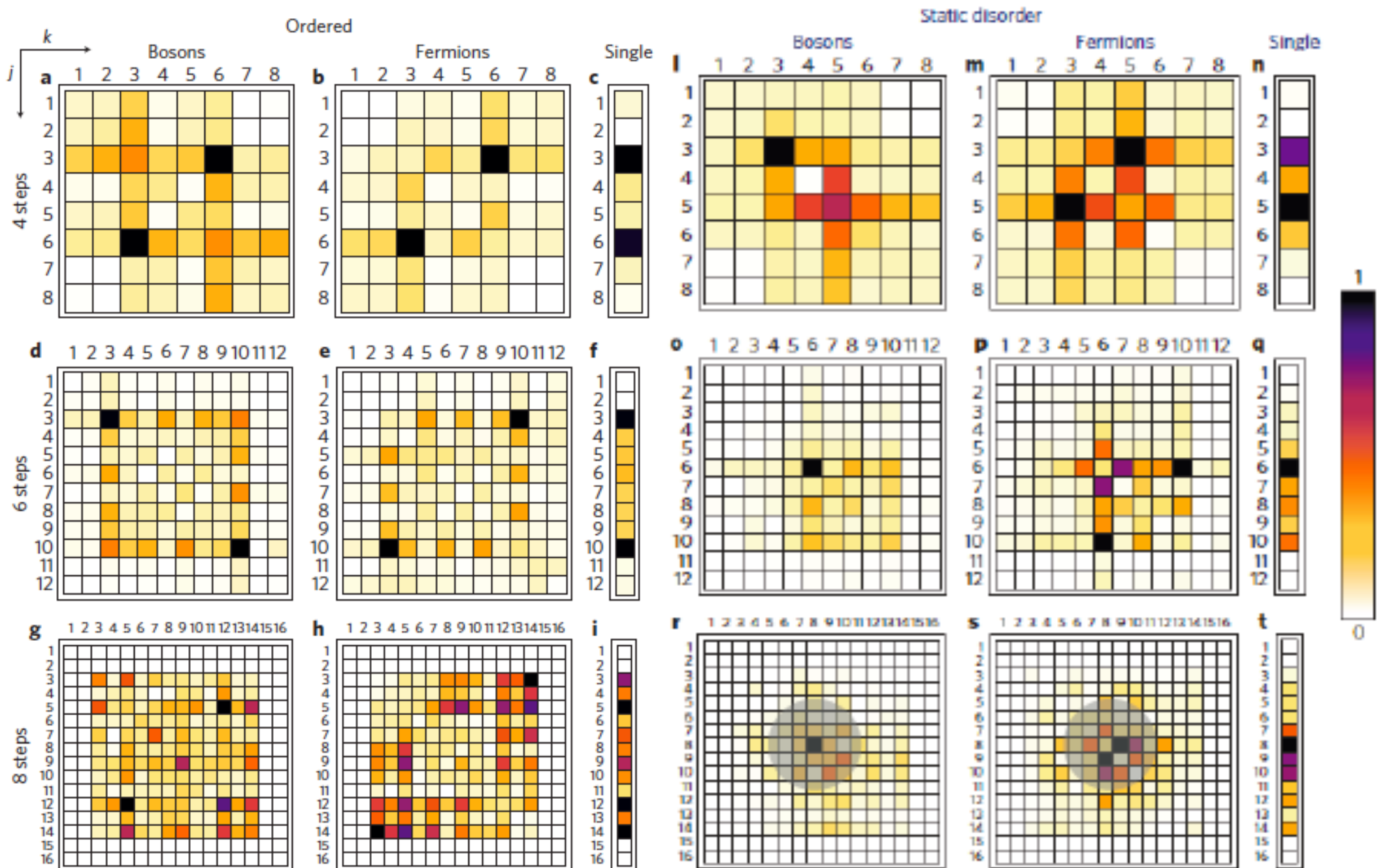
2-particle QW: variance of distributions



Ordered QW

Static disorder QW

2-particle QW: exp. results

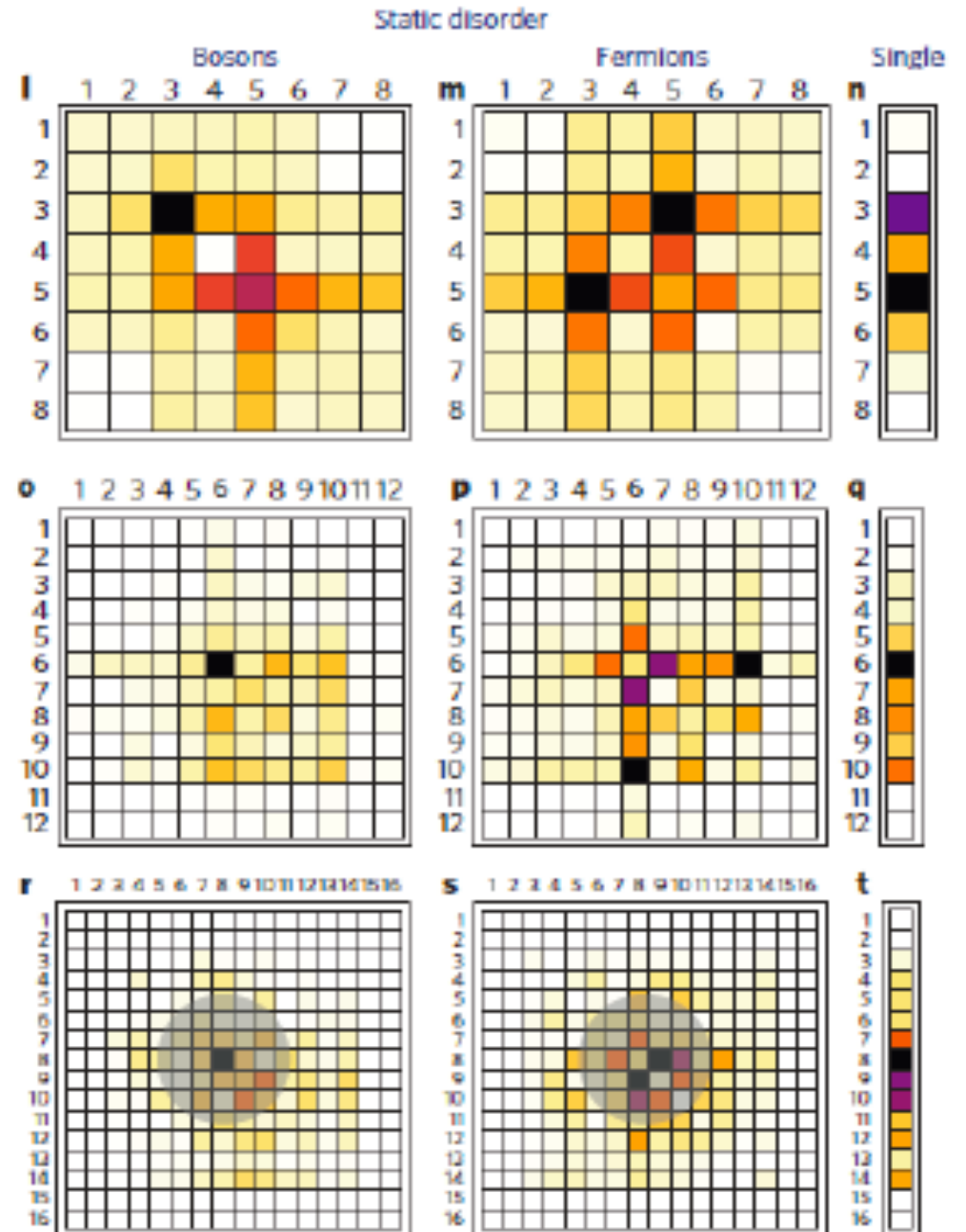


Localization evolution

4-time steps

6-time steps

8-time steps

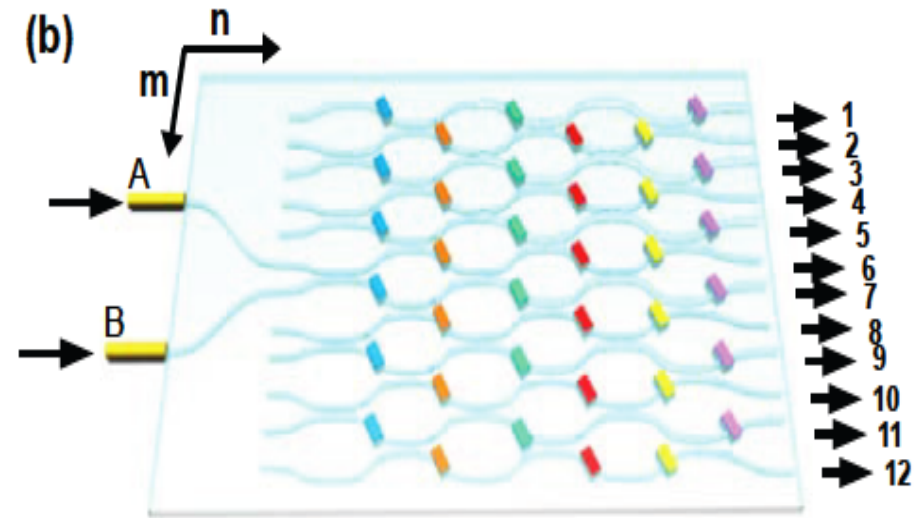


Other disordered systems

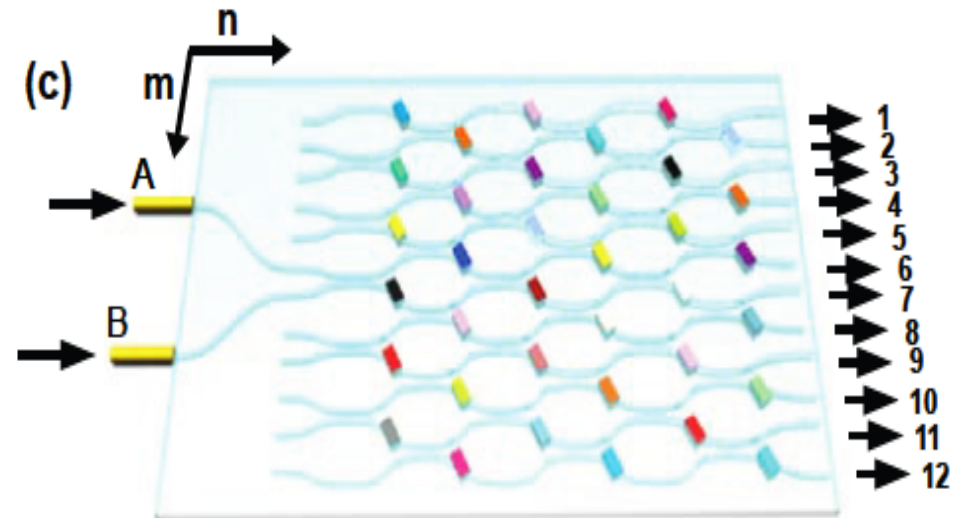
Dynamic disorder :
depending on time
but **NOT** on site



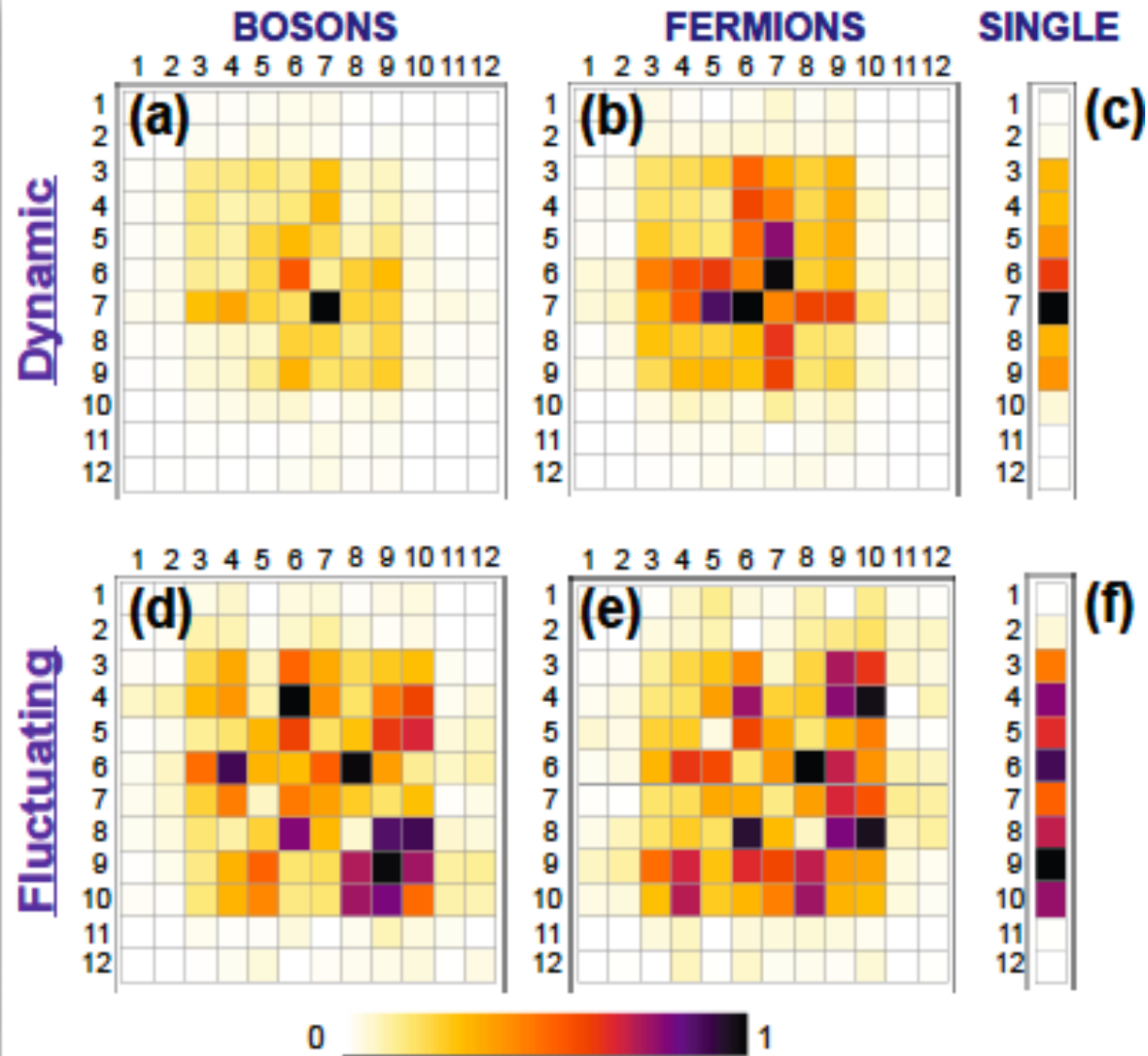
*Classical random walk
distribution*



Fluctuating disorder:
depending **BOTH**
on time and on site



Other disordered systems



The onset of dynamic disorder actually quenches the Anderson localization effects, and the distribution is far less localized. The system evolution at long times converges to a diffusion process.

Towards Quantum Supremacy: BosonSampling

HOW TO ACHIEVE QUANTUM SUPREMACY ??



John Preskill
@preskill

+ Segui

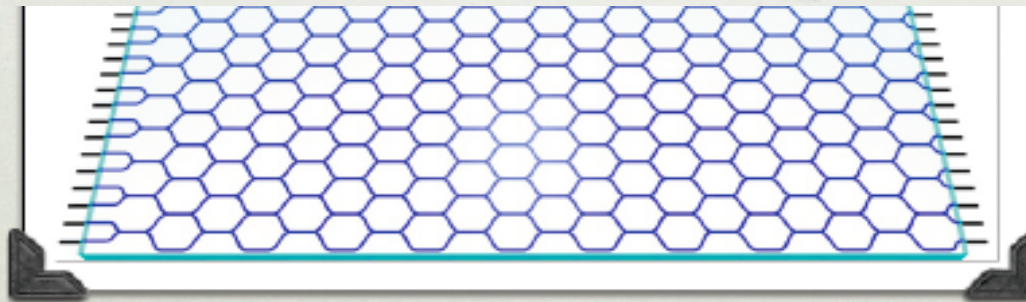
Proposed "quantum supremacy" for controlled quantum systems surpassing classical ones. Please suggest alternatives.

The Extended Church-Turing (ECT) Thesis

Everything feasibly computable in the physical world is feasibly computable by a (probabilistic) Turing machine.

Can we experimentally disproof the ECT thesis ?

Input:
 n bosons



Output:
 n -photon state

Can a classical computer simulate the distribution of the output mode numbers ?

Answer: NO!!

BosonSampling

Complex network of linear optical elements described by a $m \times m$ unitary transformation U .

Evidence of the advantage of quantum over classical computers.

Need for:

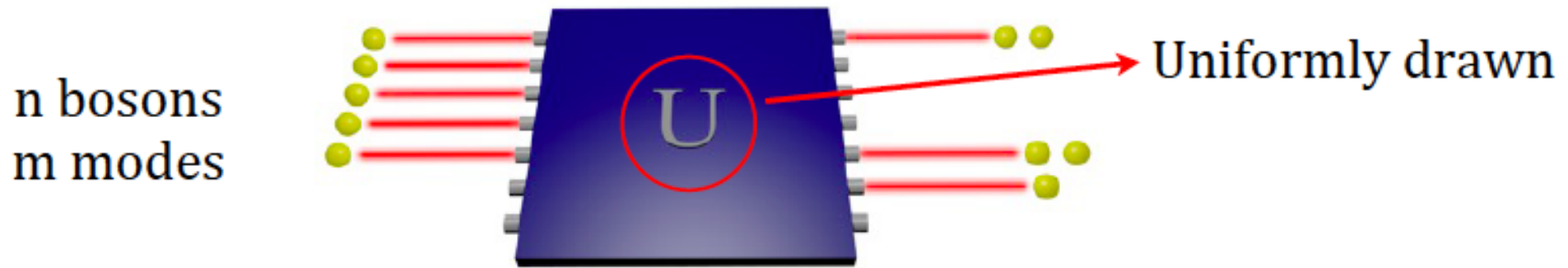
- (i) input of photons in a Fock state
- (ii) unitary evolution implemented by beam splitters and phase shifters
- (iii) simultaneous photon-counting measurement of all modes

For a large enough number of photons (10-20) and modes (100-200) classical simulation starts to be inefficient

- Permanent of the matrix: computationally hard problem (simulated by the evolution of noninteracting bosons)
- Determinant of the matrix: calculated in a polynomial time (fermion output probabilities)

BosonSampling

Sampling the output distribution (*even approximately*) of non-interacting bosons evolving through a linear network is hard to do with classical resources



Why? Transition amplitudes are related to the permanent of square matrices

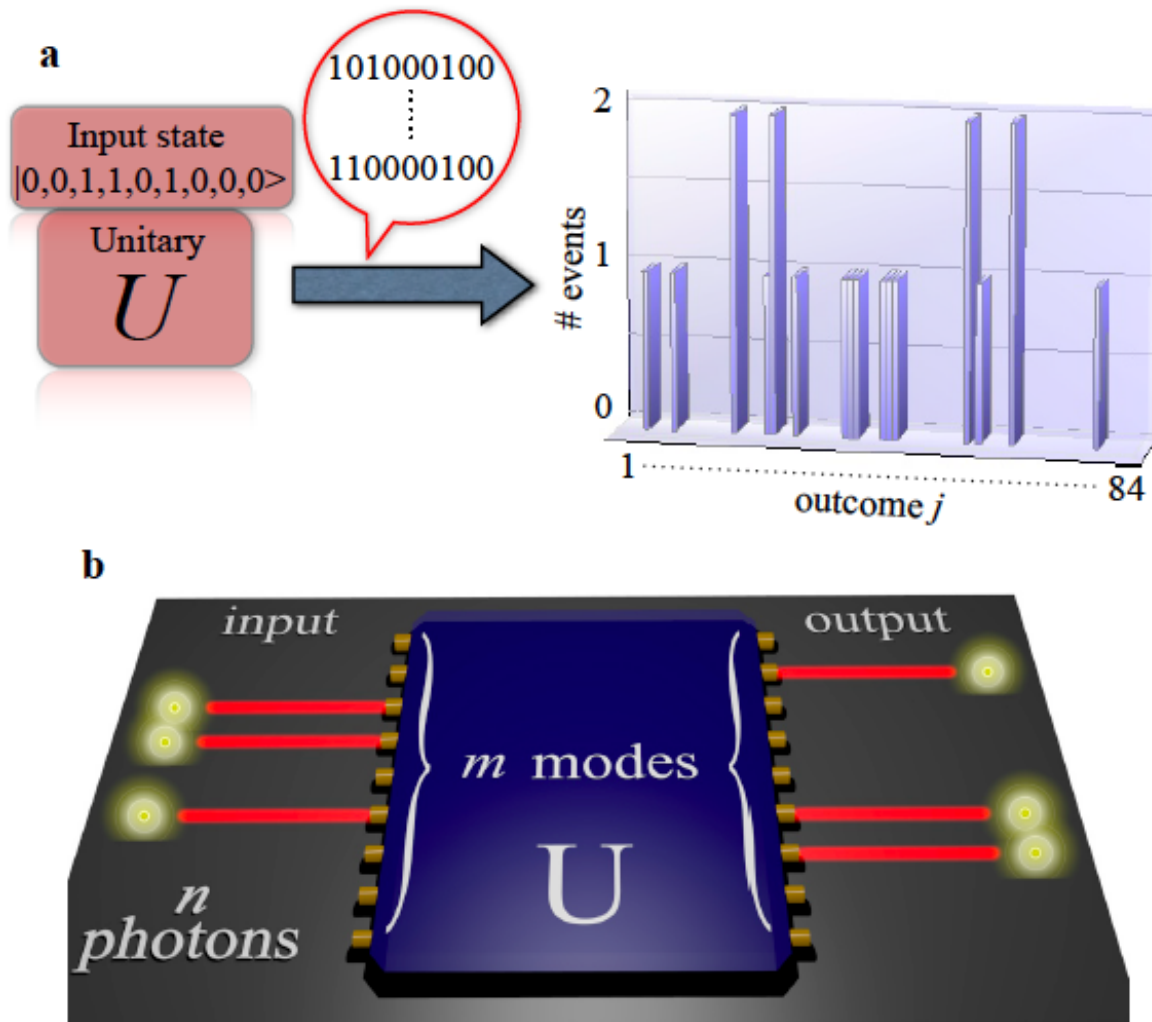
$$\langle T | U_F | S \rangle = \frac{\text{Per}(U_{S,T})}{\sqrt{s_1! \dots s_m! t_1! \dots t_m!}}$$

$$\text{Per}(A) = \sum_{\sigma \in S_n} \prod_{i=1}^n a_{i, \sigma_i}$$

classically hard

		input				
		0	1	1	0	1
output	0	0.212	-0.018 + 0.165i	-0.238 - 0.18i	-0.429 + 0.32i	-0.715 + 0.2i
	1	-0.193 - 0.388i	-0.045 - 0.379i	0.19 + 0.311i	0.328 - 0.269i	-0.594 + 0.03i
	1	-0.723 + 0.363i	0.087 - 0.09i	-0.076 - 0.155i	0.206 + 0.443i	-0.153 - 0.193i
	1	-0.092 + 0.045i	-0.148 - 0.645i	-0.588 + 0.184i	-0.369 - 0.086i	0.167 + 0.025i
	0	0.318 - 0.009i	-0.144 - 0.594i	0.452 - 0.405i	0.037 + 0.387i	0.071 + 0.025i

BosonSampling



« Small-scale quantum computers made from an array of interconnected waveguides on a glass chip can now perform a task that is considered hard to undertake on a large scale by classical means. »

Quantum optics for Boson Sampling

Photons naturally solve the Boson Sampling problem

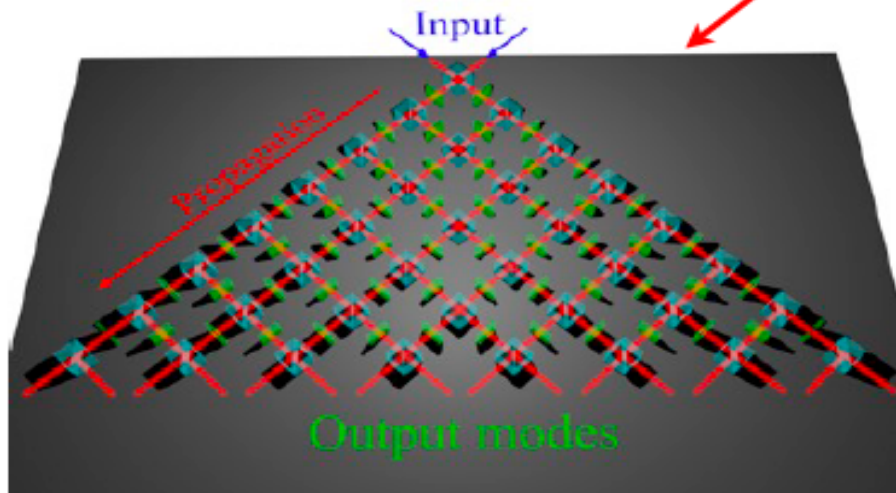
Experimental platform: photons in linear optical interferometers

Required resources: ● Single-photon inputs

● Multimode interferometers

● Detection

n photons
m modes



Hard to implement with bulk optics



Require a technological step recently available due to integrated photonics

Need:

- **High-quality multiport devices performing random unitaries with controllable characteristics**
- **Unitary reconstruction capability**

BosonSampling: an overview



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UNIVERSITY OF
OXFORD



nature
photonics

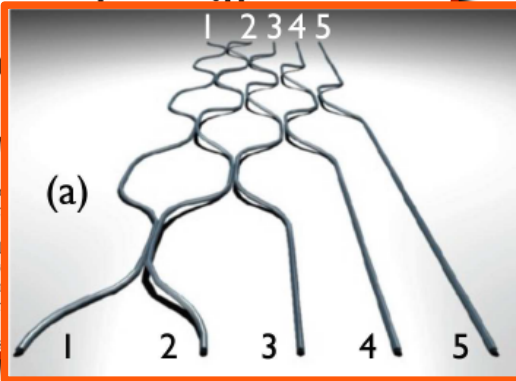
LETTERS

PUBLISHED ONLINE: XX XX 2013 | DOI: 10.1038/NPHOTON.2013.112

Integrated multimode interferometers for arbitrary designs for photonics

Andrea Crespi^{1,2}, Roberto Osellame^{1,2*}, Roberta Ramponi^{1,2}, Nicolò Spagnolo⁴, Chiara Vitelli^{4,5}, Enrico Maiorino⁴, Paolo...

1 The evolution of bosons undergoing arbitrary linear unitary transformations quickly becomes hard to predict using classical computers as we increase the number of particles and modes. 2 Photons propagating in a multiport interferometer naturally solve this so-called boson sampling problem¹, thereby motivating the development of technologies that enable precise control of multiphoton interference in large interferometers²⁻⁴. Here, we use novel three-dimensional manufacturing techniques to achieve simultaneous control of all the parameters describing...



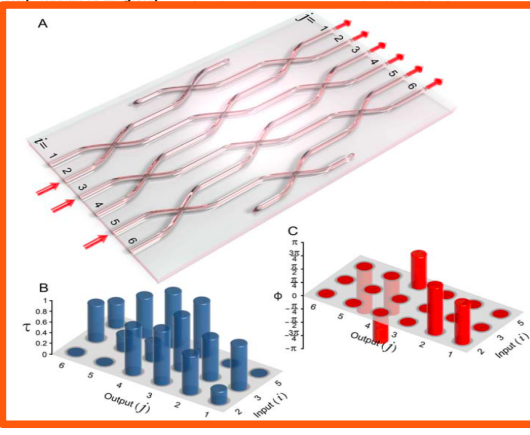
Boson Sampling on a Photonic Chip

Justin B. Spring,^{1*} Benjamin J. Metcalf,¹ Peter C. Humphreys,¹ W. Steven Kolthammer,¹ Xian-Min Jin,^{1,2} Marco Barbieri,² Animesh Datta,¹ Nicholas Thomas-Peter,¹ Nathan K. Langford,^{1,3} Dmytro Kundys,⁴ James C. Gates,⁴ Brian J. Smith,¹ Peter G. R. S...

Although universal quantum computers ideally solve problems exponentially more efficiently than classical machines, the formidable devices motivate the demonstration of simpler, problem-specific algorithms. We constructed a quantum boson-sampling machine on a photonic circuit, a problem thought to be exponentially hard to solve on a classical computer. Boson sampling merely requires indistinguishable photons, evolution, and detectors. We benchmarked our QBSM with three and four photons, demonstrating sources of sampling inaccuracy. Scaling up to larger devices could offer quantum-enhanced computation.

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modes (18). Such circuits can be rapidly reconfigured to sample from a user-defined operation (19, 20). Importantly, boson sampling requires neither nonlinearities nor on-demand entanglement, which are substantial challenges in photonics.



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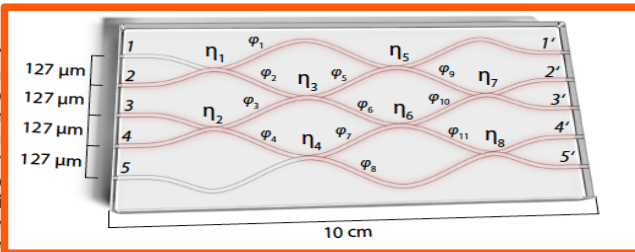
LETTERS

PUBLISHED ONLINE: 12 MAY 2013 | DOI: 10.1038/NPHOTON.2013.102

Experimental boson sampling

Max Tillmann^{1,2*}, Borivoje Dakić¹, René Heilmann³, Stefan Nolte³, Alexander Szameit³ and Philip Walther^{1,2*}

Universal quantum computers¹ promise a dramatic increase in speed over classical computers, but their full-size realization remains challenging². However, intermediate quantum computational models³⁻⁵ have been proposed that are not only hard to solve but can solve problems that are believed to be classically hard. Aaronson and Arkhipov⁶ have shown that interference of single photons in random optical networks can solve a hard problem of sampling the bosonic output distribution. Remarkably, this computation does not require measurement-based interactions^{7,8} or adaptive feed-forward techniques.



Massachusetts
Institute of
Technology



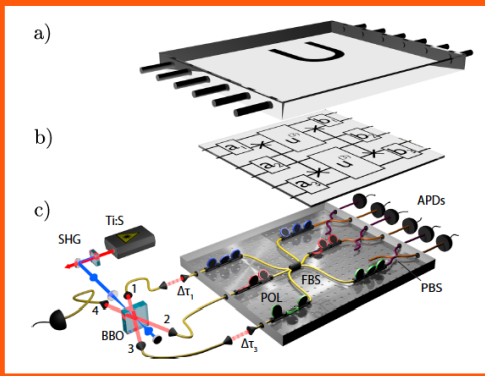
Photonic Boson Sampling in a Tunable Circuit

Matthew A. Broome,^{1,2*} Matthew J. Collins,^{1,2} Paul D. Hodgson,^{1,2} Timothy C. Ralph,^{1,2} Scott Aaronson,³ Timothy...

Quantum computers are extended Church-Turing devices that efficiently perform a sampling problem. We test a three-photon scattering unitary describing a six-photon process, even with the unavoidable scaling of a large number of modes.

To implement a circuit, the subgraphs representing circuit elements are connected by paths. Figure 4 depicts a graph corresponding to a simple two-qubit computation. Timing is important: Wave packets must meet on the vertical paths for interactions to occur. We achieve this by choosing the numbers of vertices on each of the segments in the graph appropriately, taking into account the different propagation speeds of the two wave packets [see section S4 of (32)]. In section S3.1 of (32), we present a refinement of our scheme using planar graphs with maximum degree four. By analyzing the full $(n + 1)$ -particle interacting many-body system, we prove that our algorithm performs the desired quantum computation up to an error term that can be made arbitrarily small (32). Our analysis goes beyond the scattering theory discussion presented above; we take into account the fact that both the wave packets and the graphs are finite. Specifically, we prove that by choosing the size of the wave packets, the number of vertices in the graph, and the total...

15 FEBRUARY 2013 VOL 339



Paper	Group	Contents	Validation
Science 339, 794 (2013)	Brisbane, Boston	n=2,3 photons, m=6 modes - fiber optics	No
Science 339, 798 (2013)	Oxford	n=3 photons, m=6 modes + n=4 photons with (lower complexity) bunched input	No
Nat. Photon. 7, 540 (2013)	Vienna, Jena	n=3 photons, m=5 modes	No
Nat. Photon. 7, 548 (2013)	Roma, Milano, Niteroi	n=3 photons, m=5 modes Haar-Random unitary	No
PRL 111, 130503 (2013)	Roma, Milano, Niteroi	Bosonic Birthday paradox, and verification of full-bunching law	No
Nat. Photon. 8, 615 (2014)	Roma, Milano, Niteroi	n=3 photons, m=5,7,9,13 modes validation tests	Uniform distribution, distinguishable particles
Nat. Photon. 8, 621 (2014)	Bristol	n=3 + n=4,5 photons (subtracting bunching), m=21 qwalk n=3 photons in m=9 Haar Unitary	Uniform distribution, distinguishable particles
Phys. Rev X, 5, 041015 (2015)	Vienna, Jena	investigation on complexity with partial photon distinguishability, n=3 photons, m=5 modes	No
Science Advances 1, e1400255 (2015)	Roma, Milano, Niteroi	n=3 photons, m=9,13 modes scattershot of 8 input states	Uniform distribution, distinguishable particles
Science 349, 711 (2015)	Bristol	implementation of 6x6 fully reconfigurable circuit, Haar-random. n=3: zero-transmission in Fourier matrix n=6 with bunched input (2 modes)	Distinguishable particles
Nat. Commun. 7, 10469 (2016)	Roma, Milano	n=2 photons, m=4,8 modes suppression law in Fourier matrix with scalable 3D architecture	Distinguishable particles, mean-field state

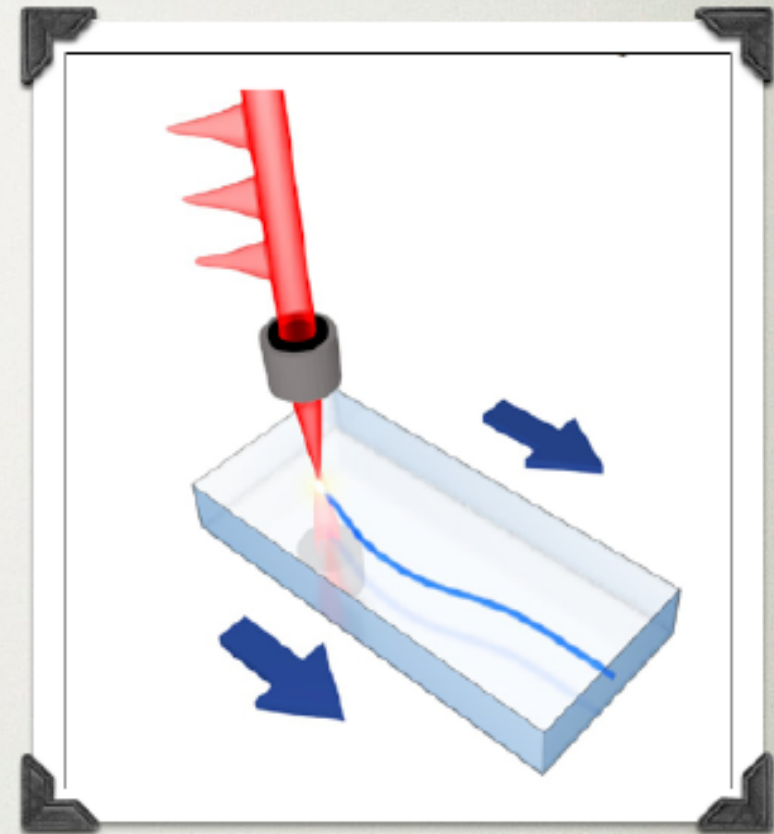
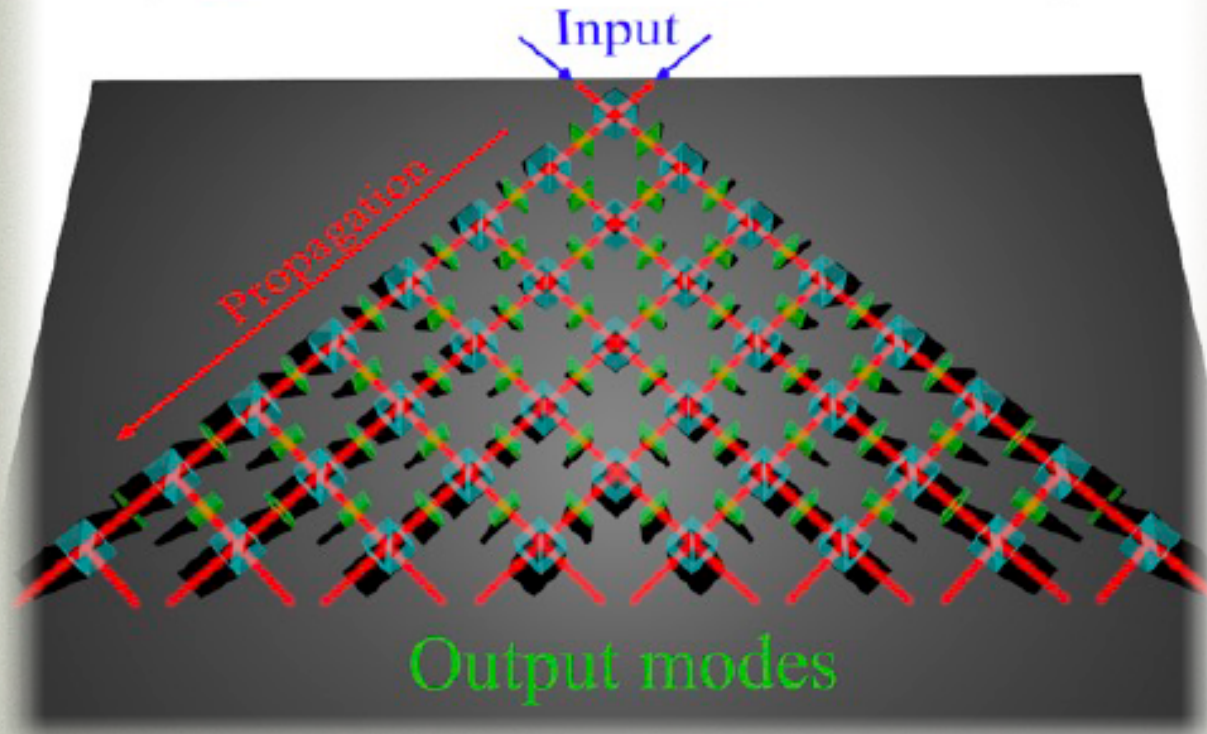
THE SOLUTION: INTEGRATED PHOTONICS



beam-splitter



phase shift



Laser writing technology:

unique capability to transmit any polarization state

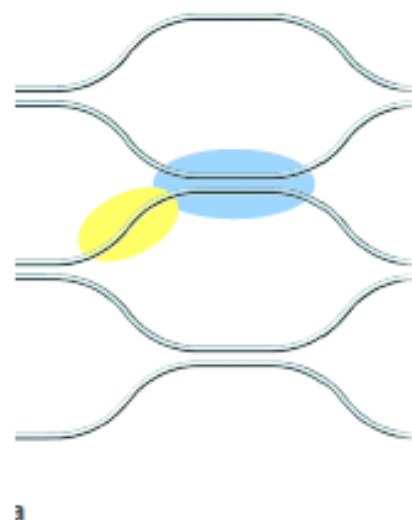
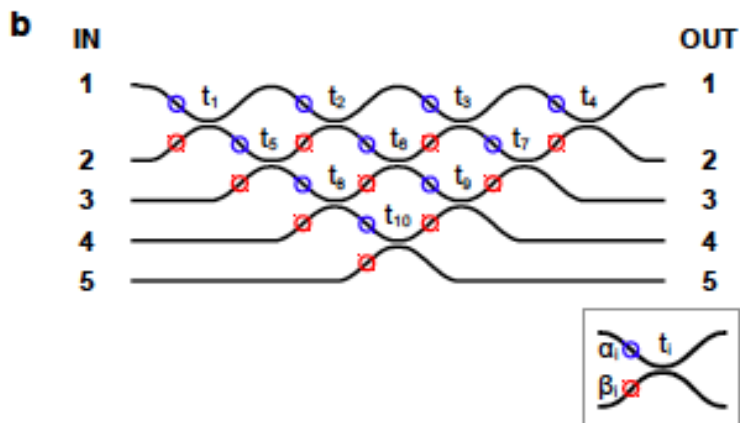
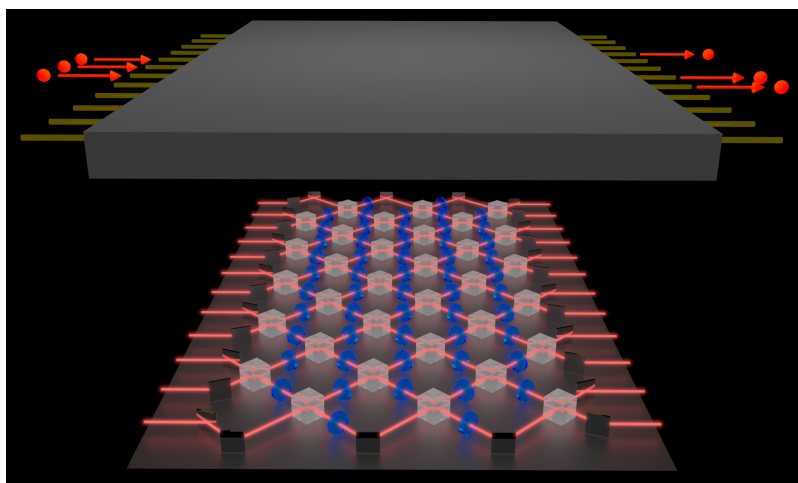
- Femtosecond pulse tightly focused in a glass
- Waveguides writing by translation of the sample

Our approach: controlling ϕ and T

Requirement for Boson Sampling -
design arbitrary interferometers

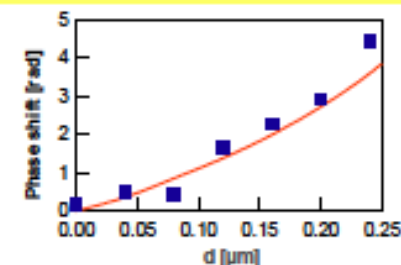
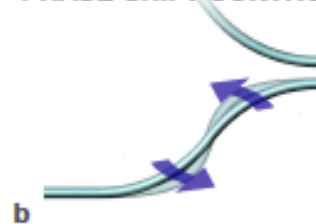


Requires independent control of
phases and beam-splitter operation

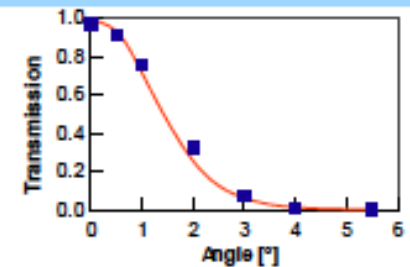


Fabrication process

PHASE-SHIFT CONTROL



TRANSMISSION CONTROL

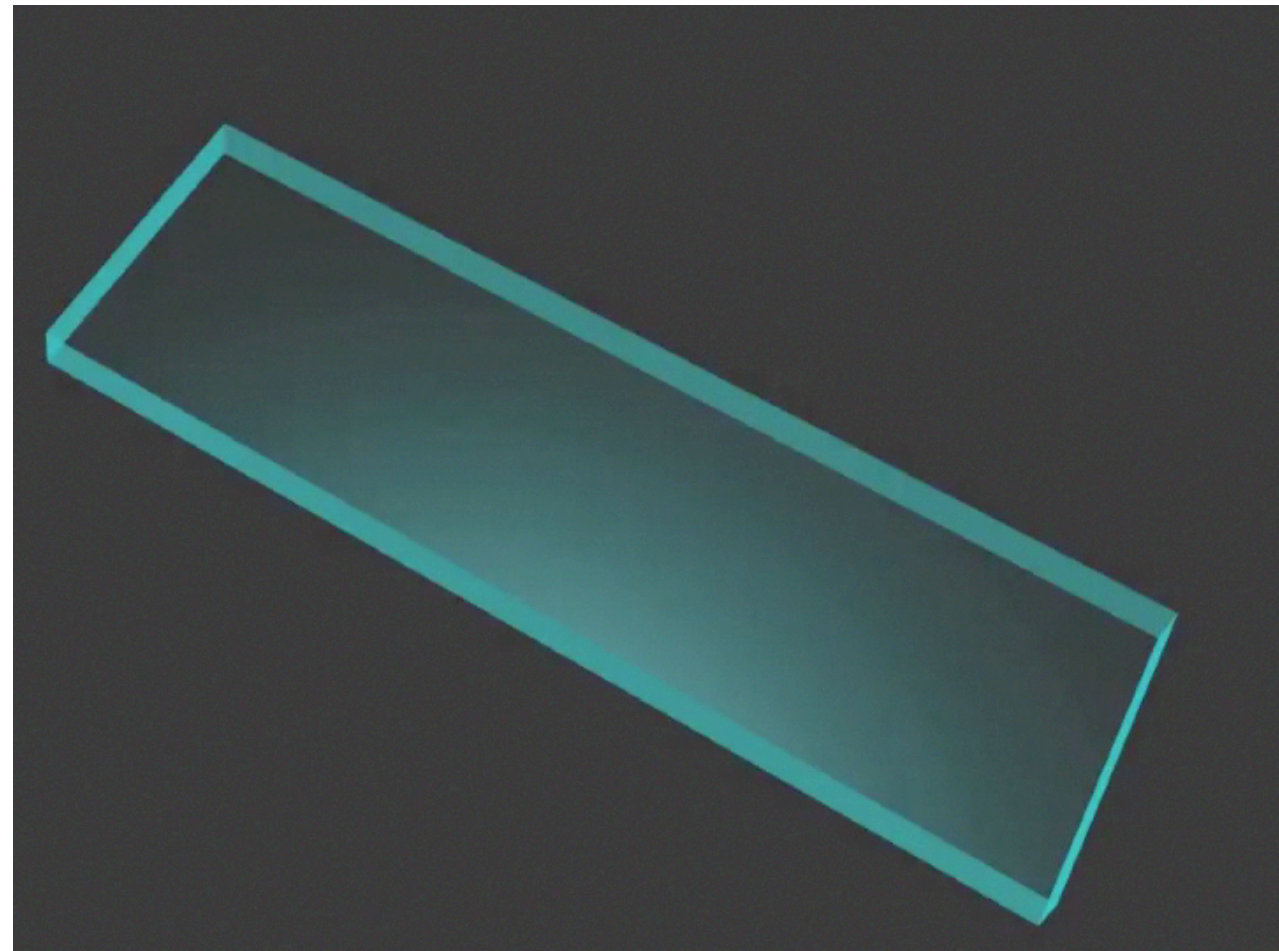
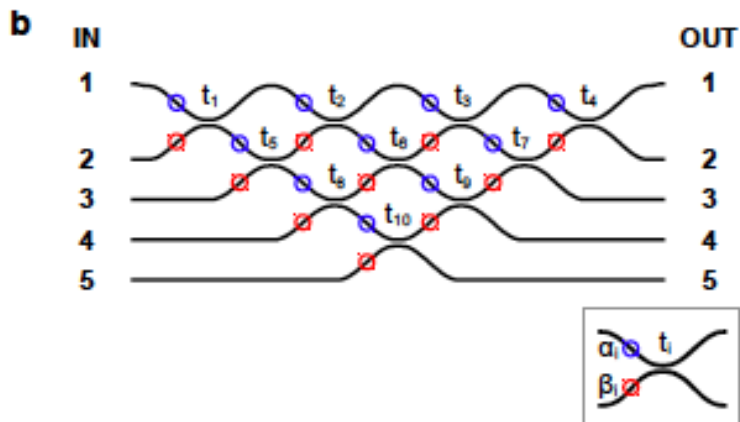
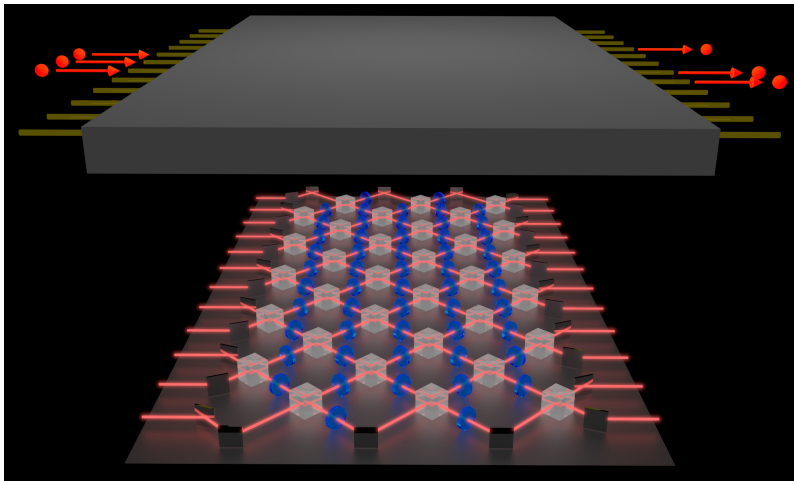


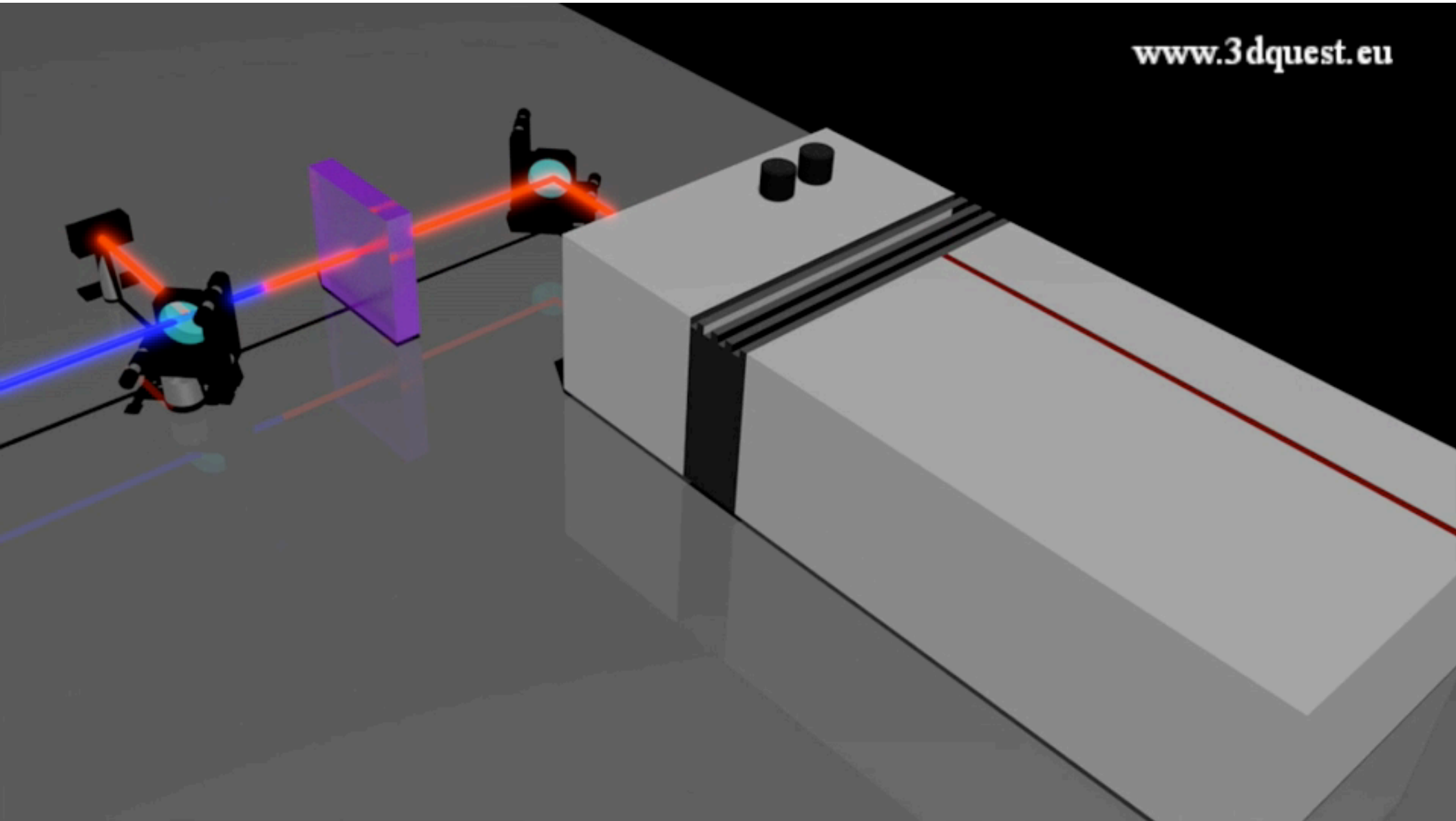
The chip

Requirement for Boson Sampling -
design arbitrary interferometers



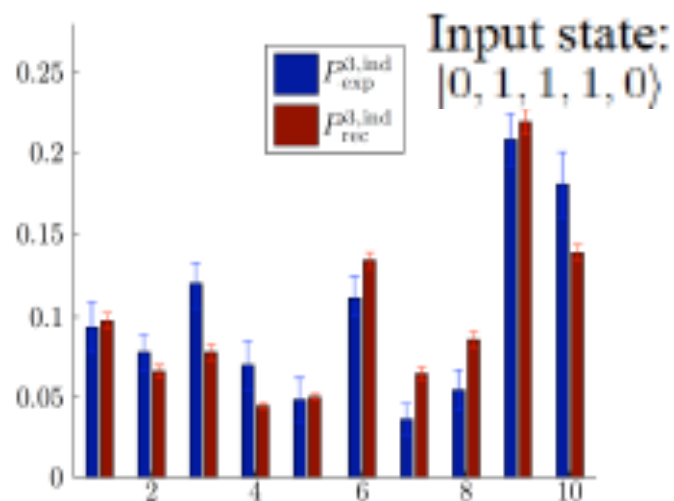
Requires independent control of
phases and beam-splitter operation





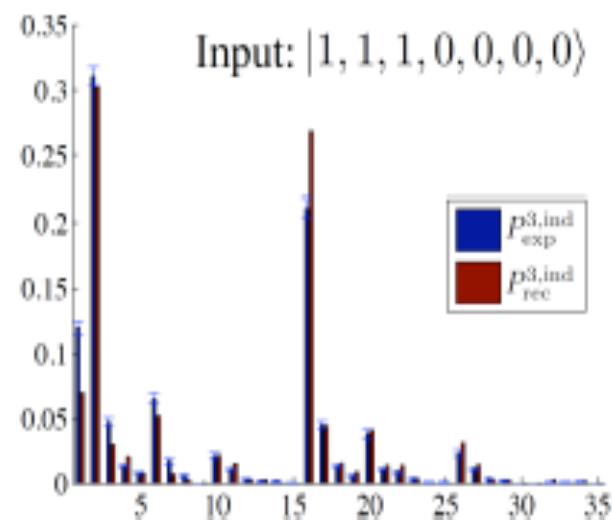
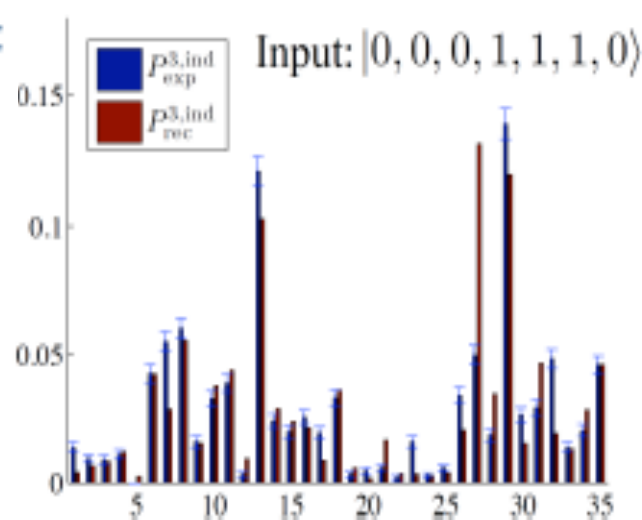
3 photons, 5 and 7 modes

a 5-modes interferometer



b

7-modes interferometer



$$D_{exp, sp} = 0.105 \pm 0.024$$

$$D_{exp, sp} = 0.122 \pm 0.037$$

Partial distinguishability of the photons taken into account

- Confirmation of Permanent formula
- Upgraded to a higher number of modes

Can Boson Sampling be validated?

It has been argued that due to the high complexity, BosonSampling output in the hard-computational regime cannot be distinguished from the random output of a uniform distribution

C. Gogolin et al. *arXiv:1306.3995*

The Theorists' Answer

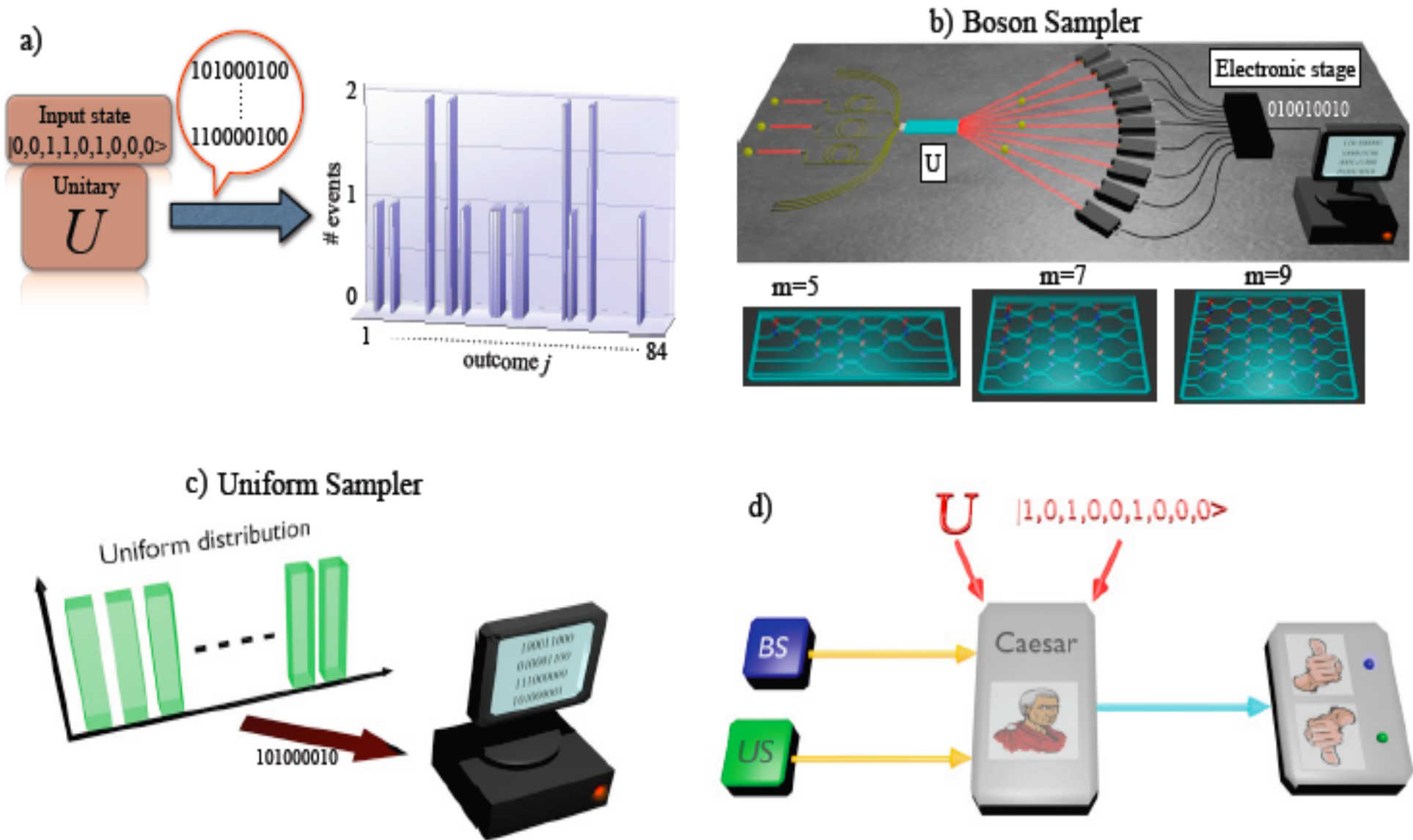
For each single registered event, which identifies the output state, calculate the quantity

$$P = \prod_{i=1}^n \sum_{j=1}^n |A_{i,j}|^2$$

whith $A_{i,j}$ = submatrix of U depeing on the input and output states, and compare this value to its counterpart for a uniform distribution P_u .

If $P > P_u$, you can guess that the single event has been produced by a BosonSampler

Validation of BosonSampling

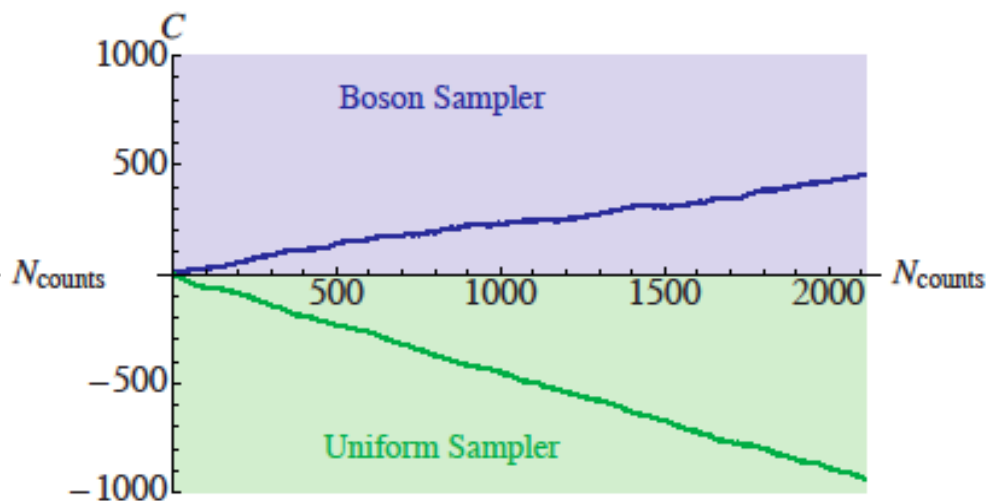
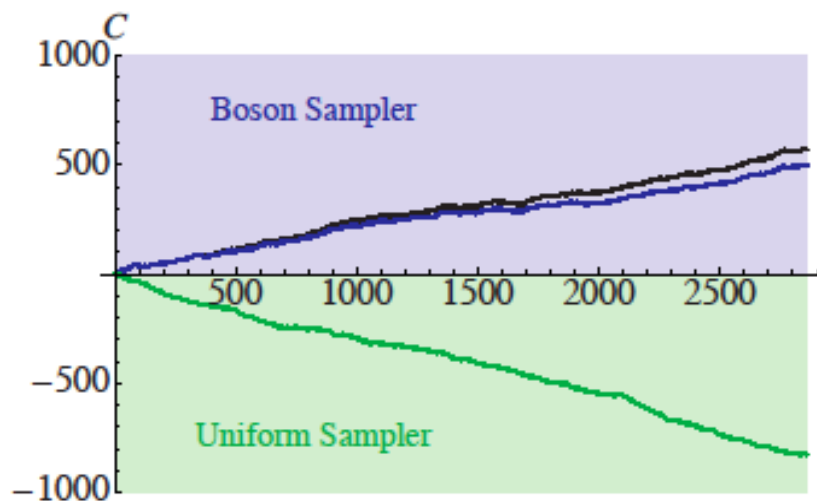
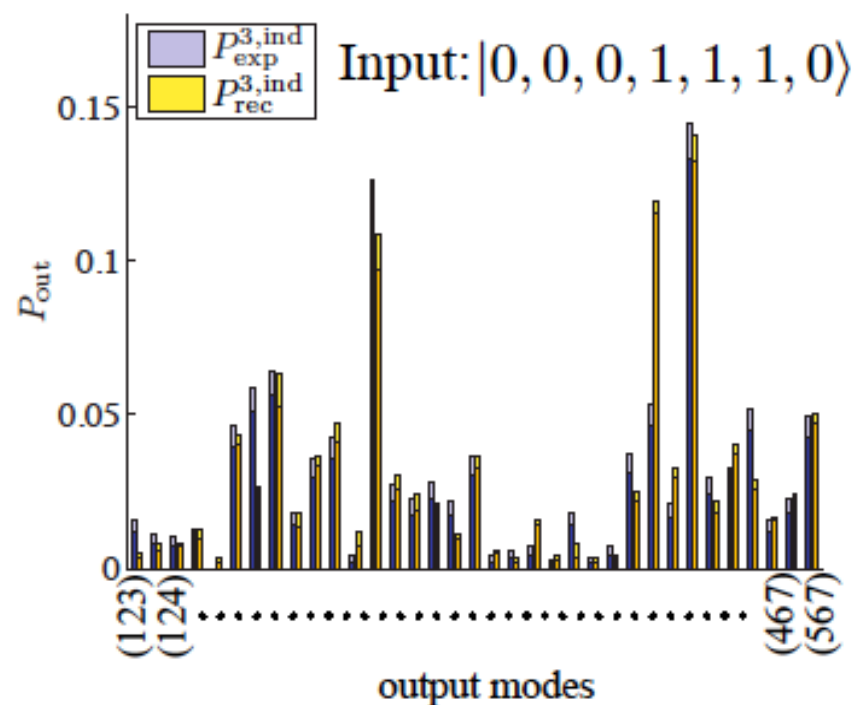
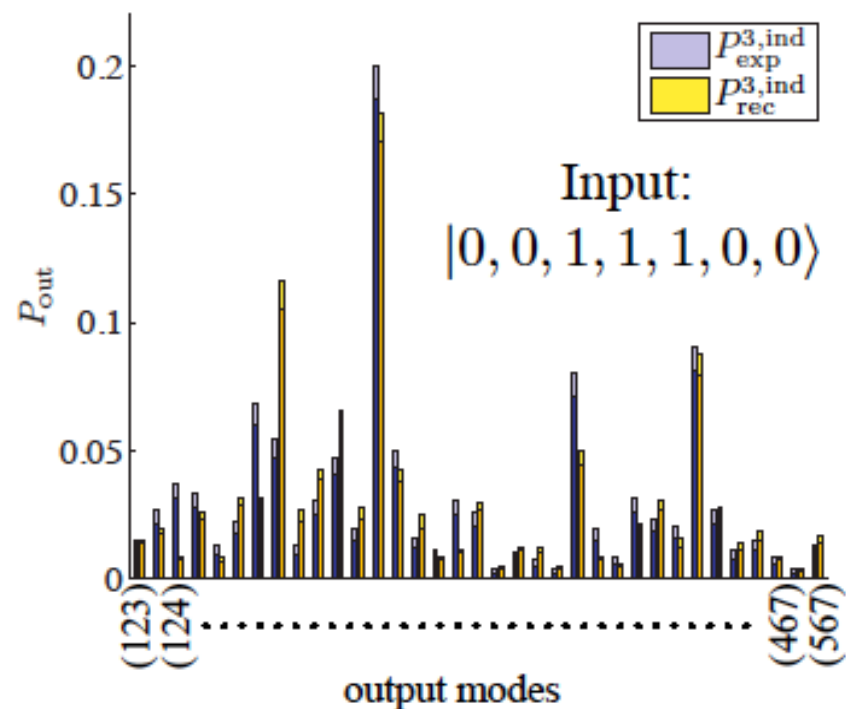


N. Spagnolo, et al., *Nature Photonics* **8**, 614 (2014)

Similar experiment in Bristol: J. Carolan, et al., *Nature Photonics* **8**, 619 (2014)

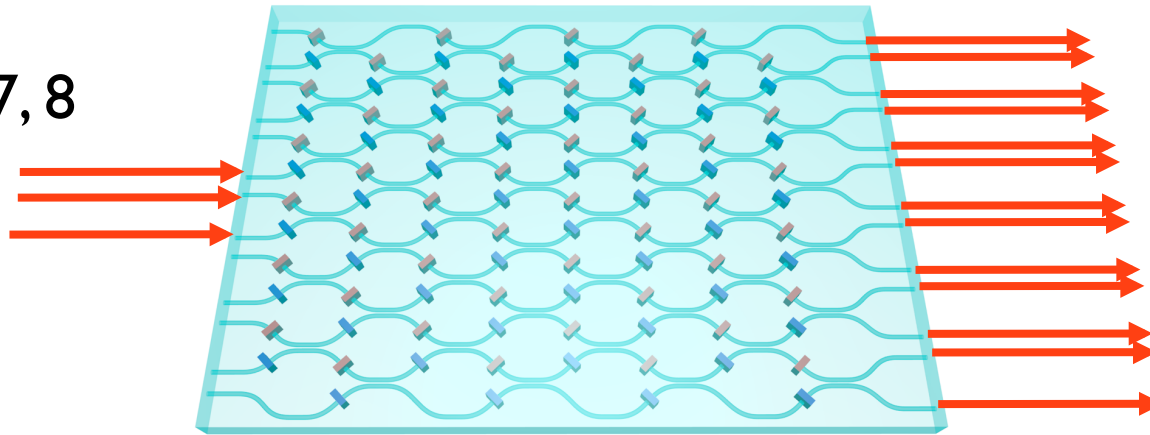
Experimental proof of validation, 7 modes

b) 7-mode interferometer

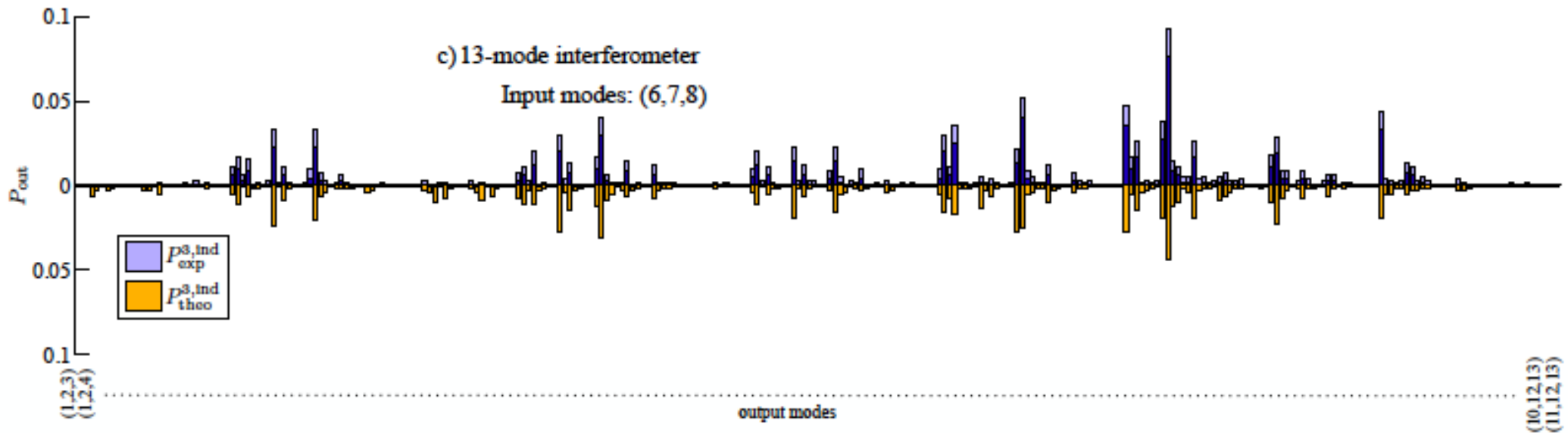


BosonSampling in a 13 mode system

Input: 6, 7, 8



Output: 286 different possible no-bunching configurations



Towards Quantum Supremacy

GOAL: Achieve Boson Sampling with $n = 10-20$ photons and $m = 100-200$ modes

Open questions

- *Measure BS complexity*
- *Other equivalent experimental schemes*
- *Certify the functioning of a BS experiment*
- *How noise/imperfections affect a complex BS*

Challenges

- *Efficient single photon sources*
- *Reconfigurable photonic circuits*
- *Efficient single photon detectors*

Scattershot Boson Sampling

p = probability of generating a photon pair in a single source
(typical values $p=0.01-0.015$)

p^n probability of generating the n -photon input

Scattershot Boson Sampling, n -photon term

$p^n (1 - p)^{m-n}$ probability of generating one of the n -photon input configurations

$\binom{m}{n}$ number of possible output configurations

Total generation rate: $\sim p^n (1 - p)^{m-n} \binom{m}{n}$

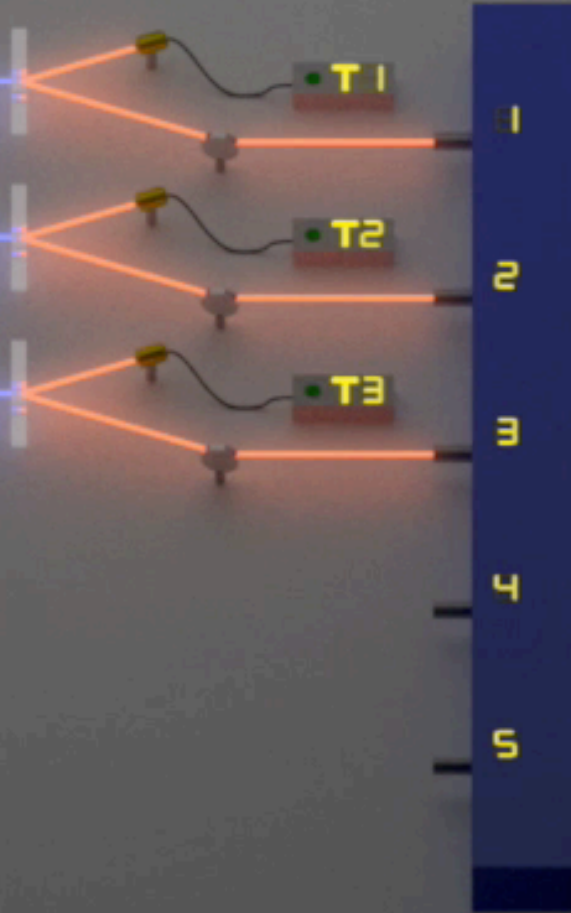
Sample both from the *input* and the *output modes*



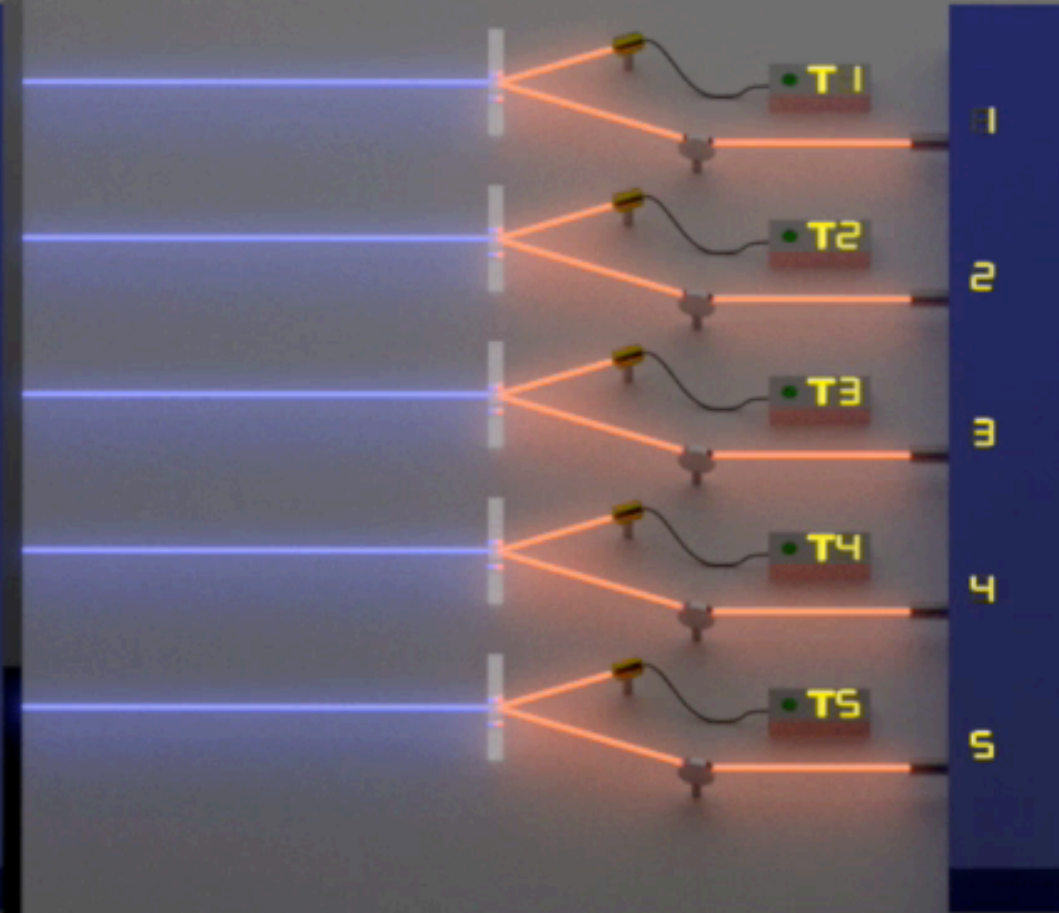
Potentially huge increase of the brightness of the quantum hardware

Scattershot BosonSampling

BOSON SAMPLING

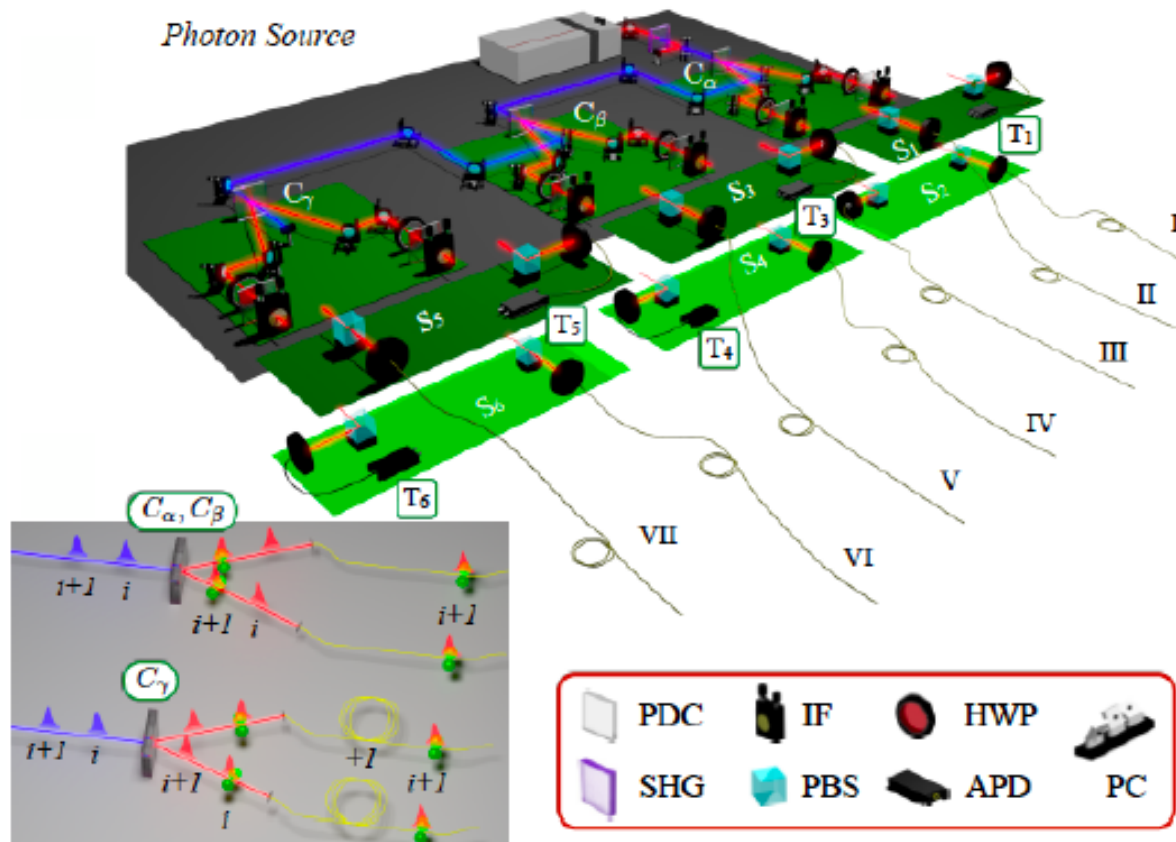
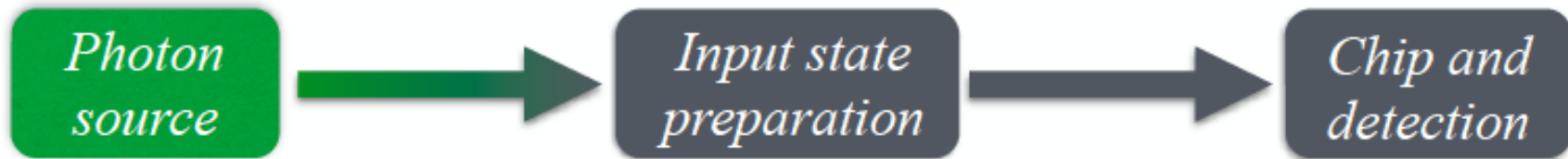


SCATTERSHOT BOSON SAMPLING



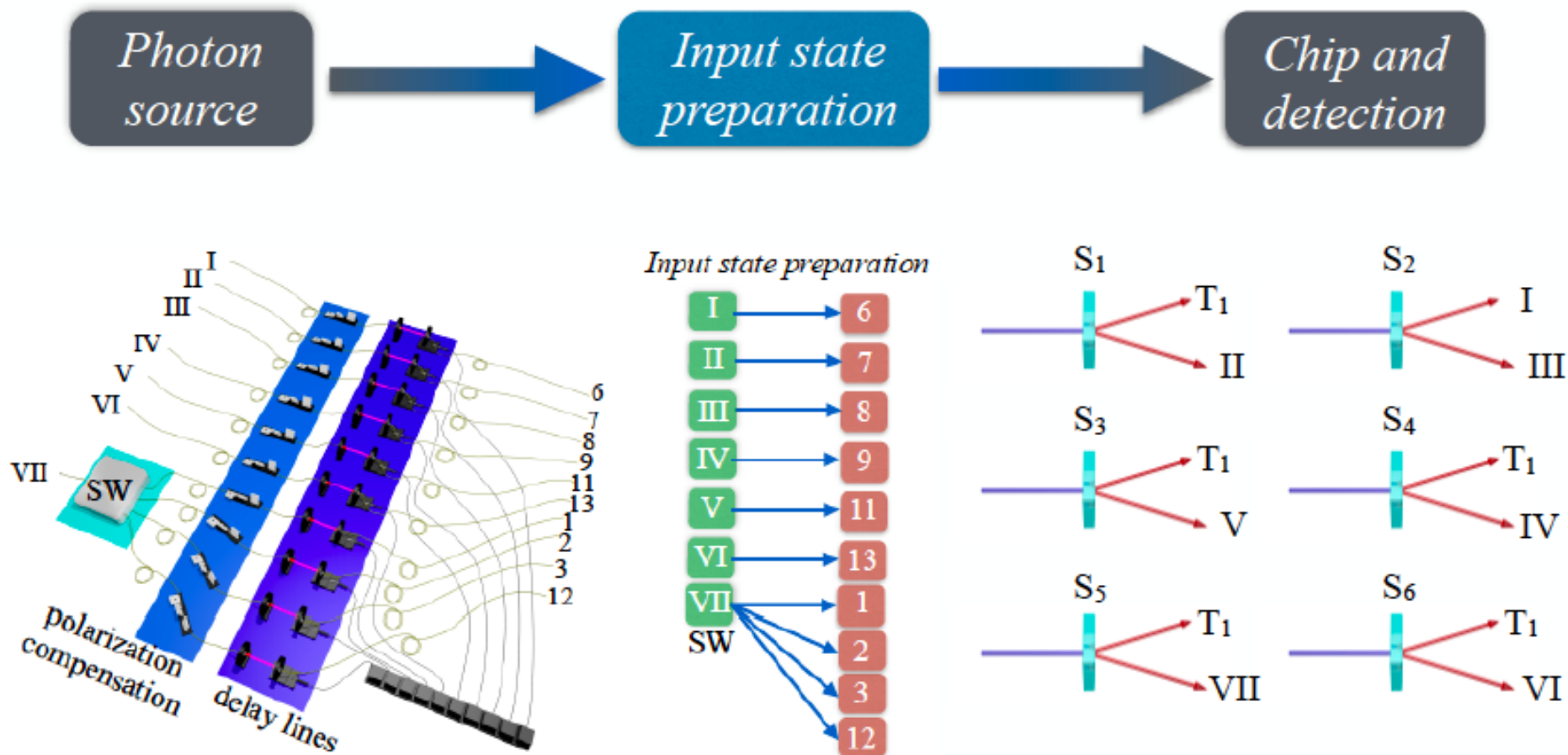
Scattershot BosonSampling: generation

Experimental setup - 1



Scattershot BosonSampling: preparation

Experimental setup - 2

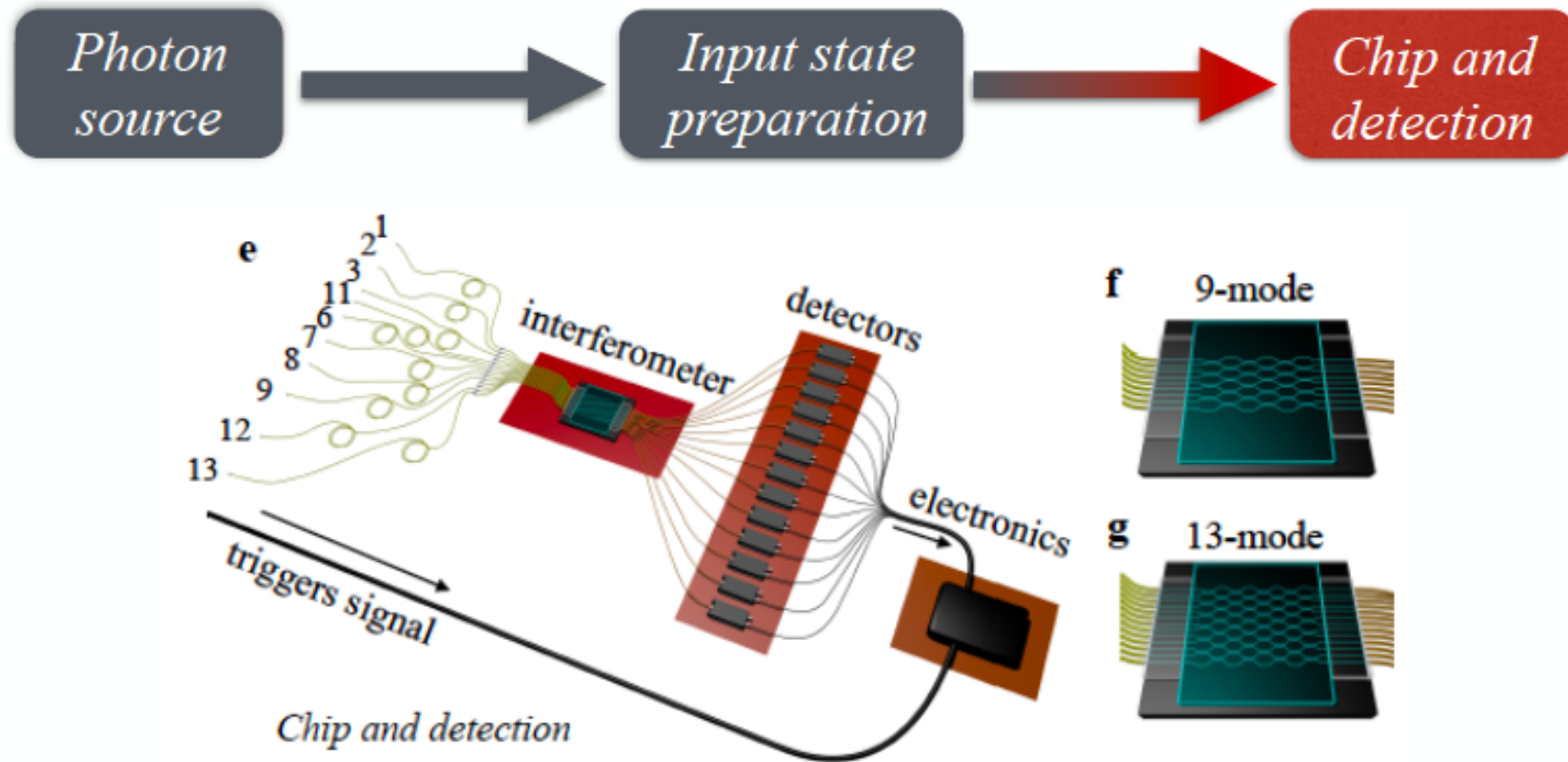


- Three photon events:
- 1) Photon I (input 6) [fixed]
 - 2) Photon III (input 8) [fixed]
 - 3) Random input heralded by T_i

Input randomness further enhanced by sequential switching of photon VII

Scattershot BS: chip and detection

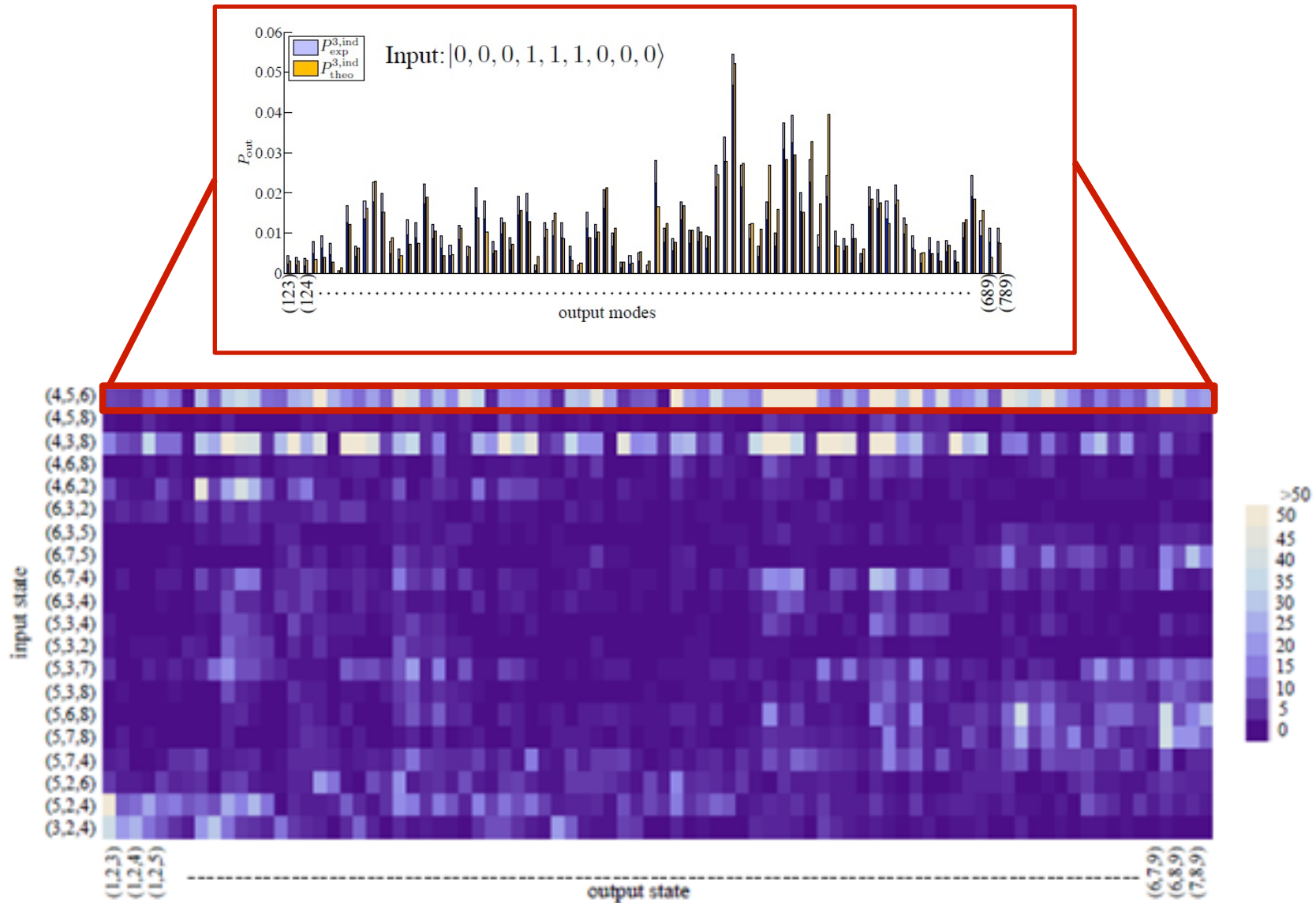
Experimental setup - 3



Evolution through $m=9$ and $m=13$ interferometers with random (but known) structure

Coincidence detection for:
Three-photon events with one heralding trigger
Two-photon events with two heralding triggers

Scattershot BS: random input



Certification of many-body interference

Boson Sampling experiments pose serious problem of certification of the result's correctness in the computationally-hard regime.

Use 3-D photonic chips to test true n-photon interference in a multimode device [by M.C. Tichy *et al.* (Phys. Rev. Lett., 2014)].

Proposal based on the suppression of specific output configurations in an interferometer implementing an n^D -dimensional Quantum Fourier Transform (QFT) matrix.

Generalization of the 2-photon/2-modes Hong-Ou-Mandel (HOM) effect, used to test a wide range of photonic platforms.

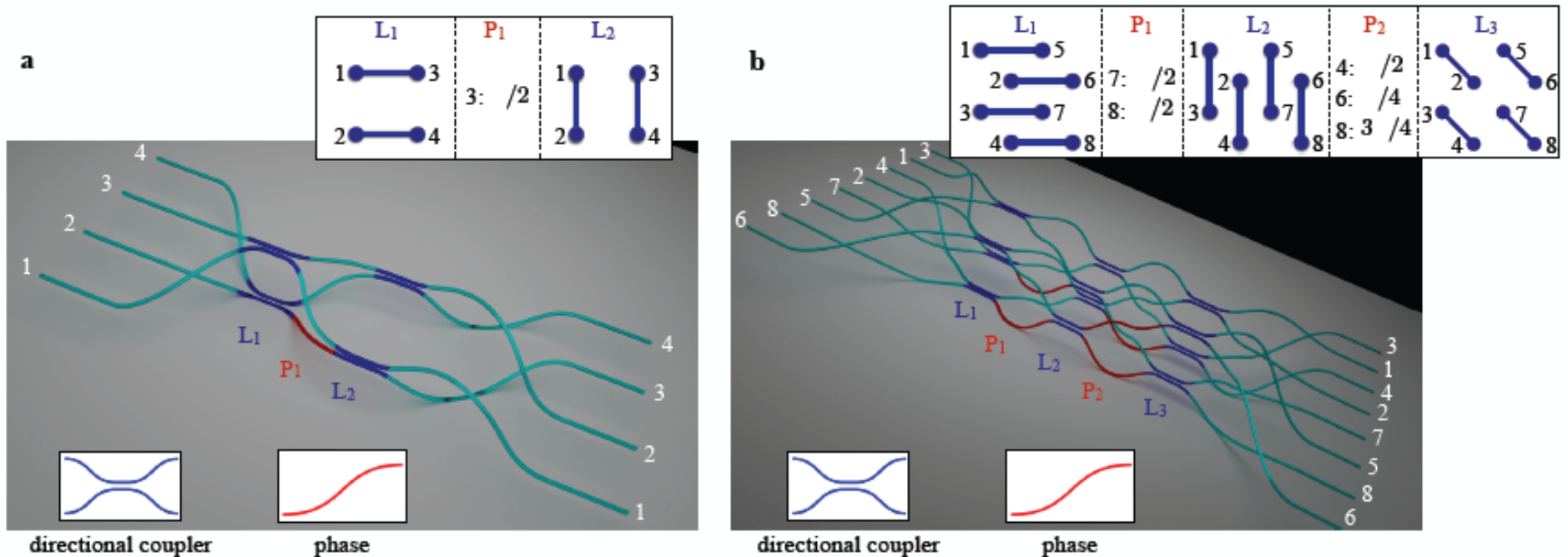
Quantum Fourier Transform in a 3D chip

Quantum interference in multimode interferometers may determine suppression of a large fraction of the output configurations.

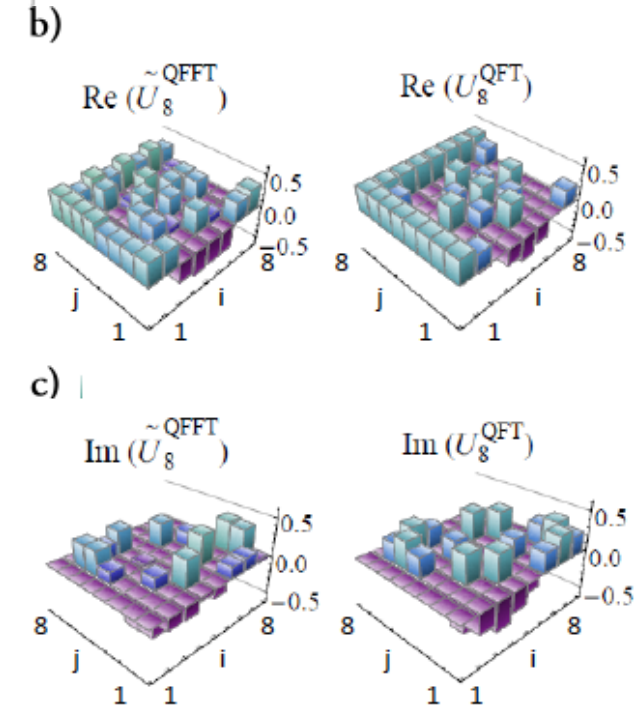
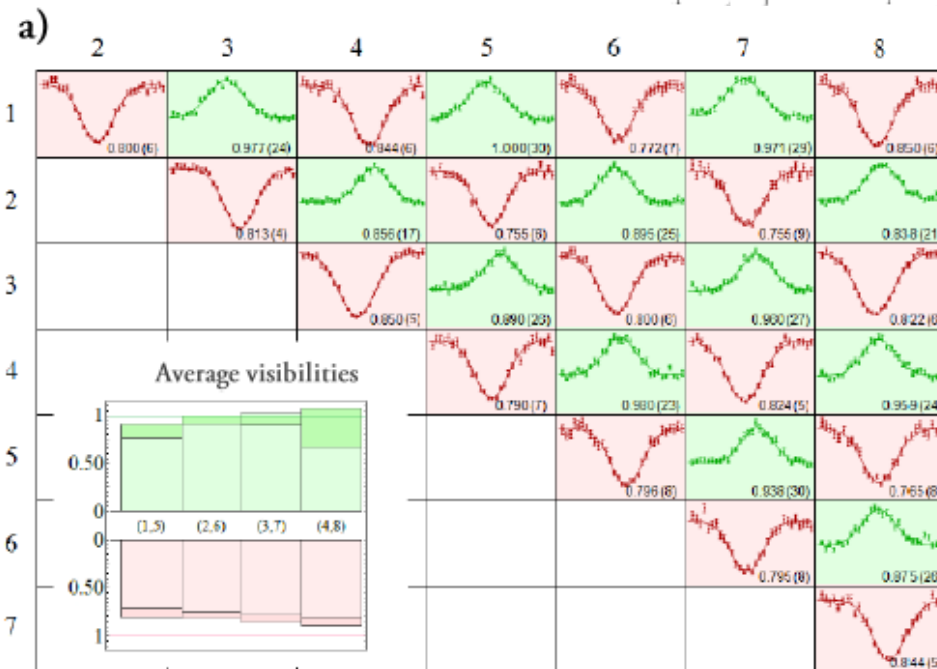
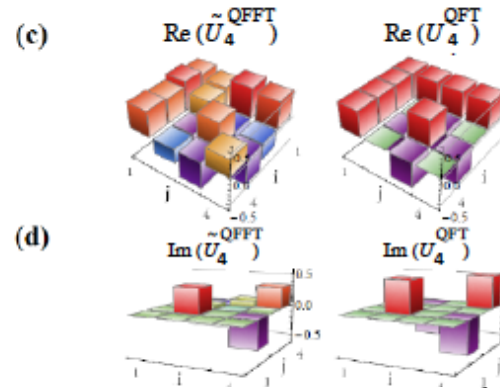
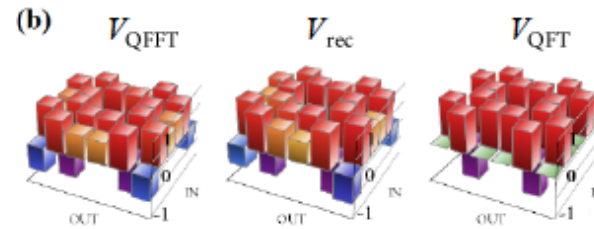
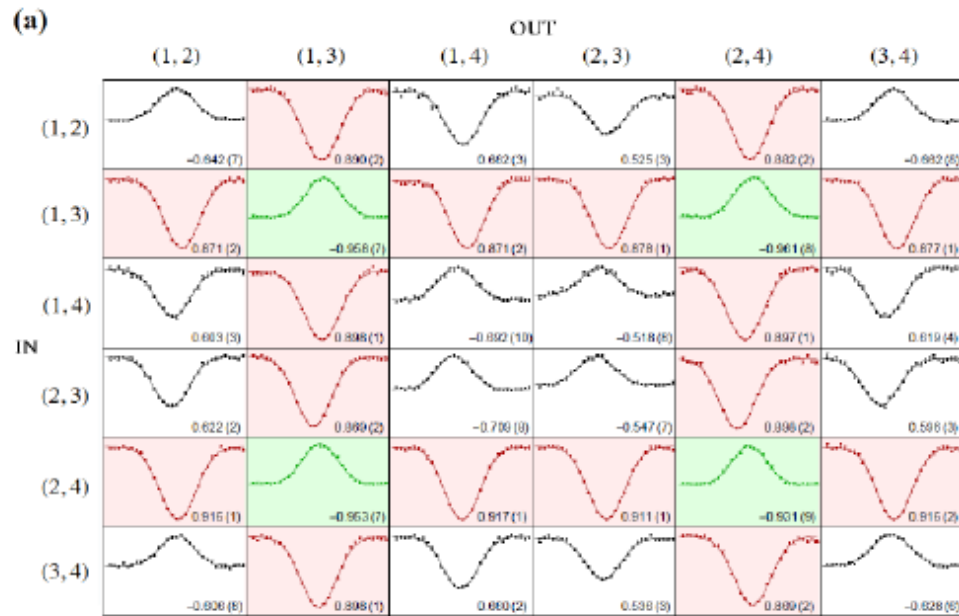
Study the evolution of particular input states through the network implementing the QFT described by the unitary matrix:

$$U_{l,q}^{\text{QFT}} = \frac{1}{\sqrt{m}} e^{i \frac{2\pi lq}{m}}$$

Test performed with 2 photon and 4- and 8- mode interferometers



Quantum suppression law in a 3D chip



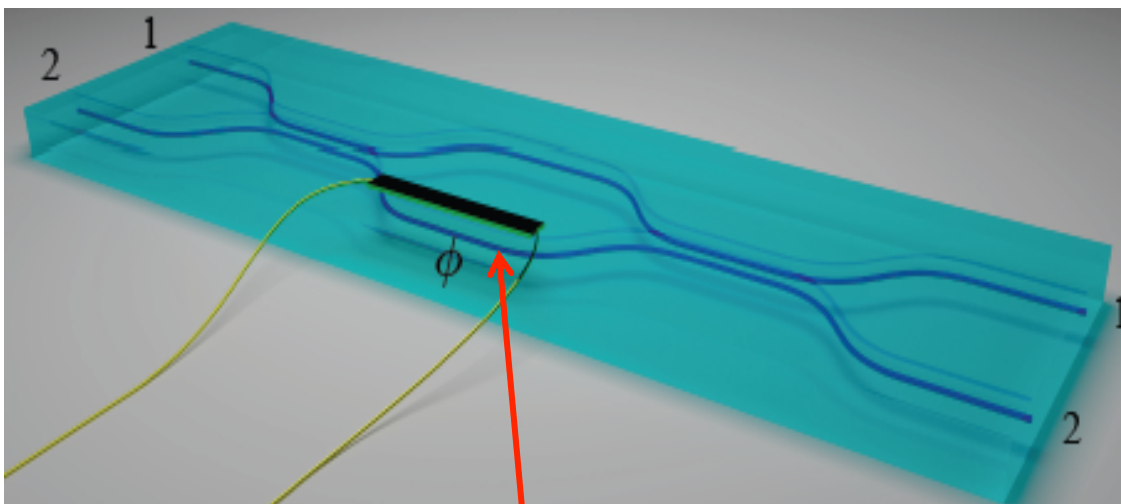
Building blocks of integrated photonics

- Directional coupler (BS)
- Mach – Zehnder interferometer
- Phase shifter
- Polarizing directional coupler (PBS)
- Waveplate (HWP, QWP)

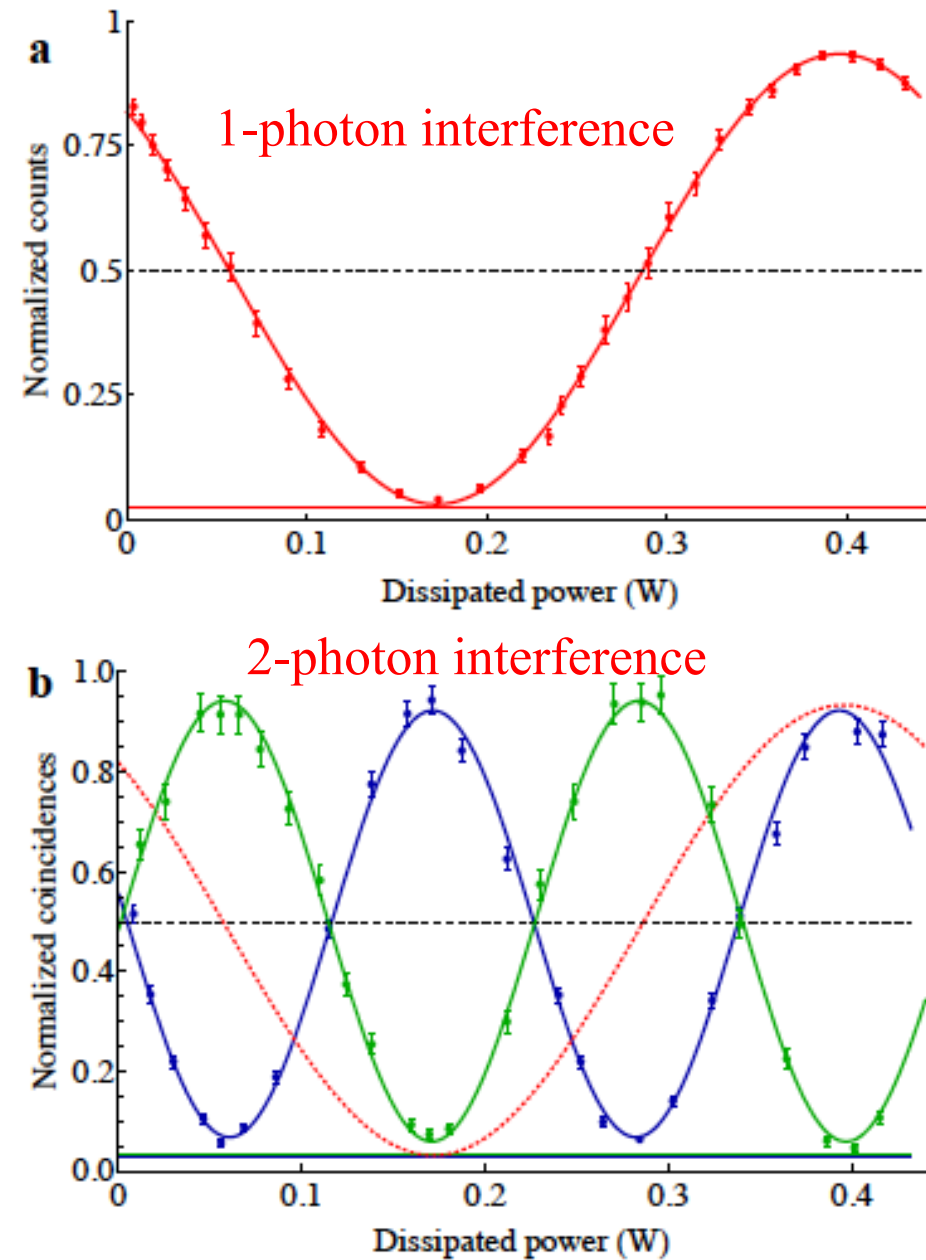
Use with single- or multi-photon states, with path and/or polarization qubit encoding

Two-arm interferometer – Phase shifter

Mach-Zehnder interferometer



Phase ϕ changed by electric thermistor

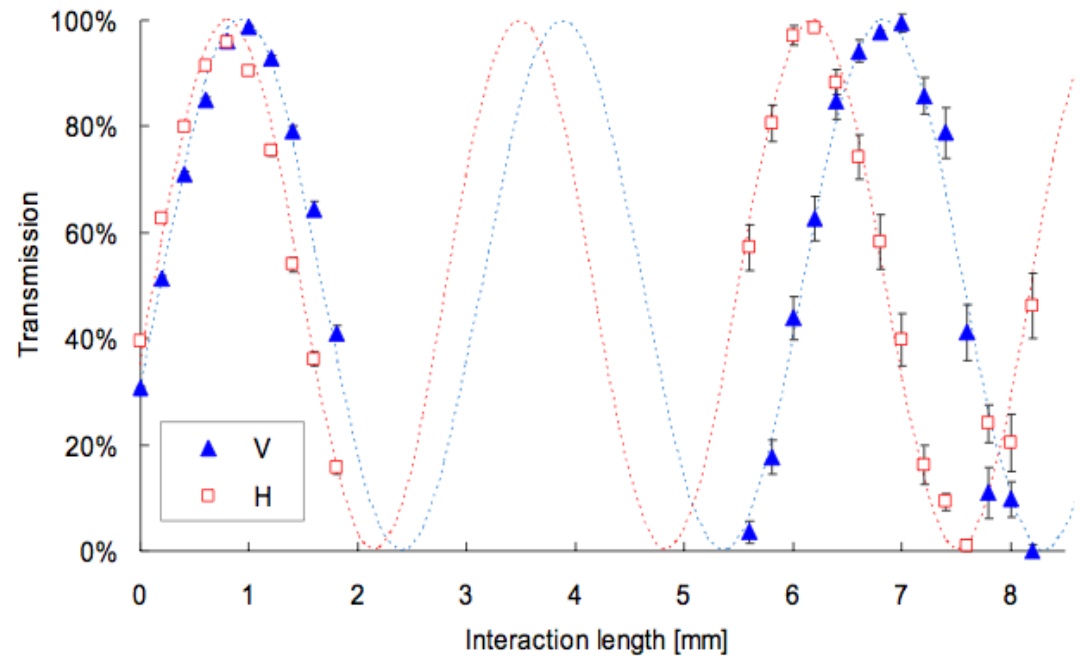
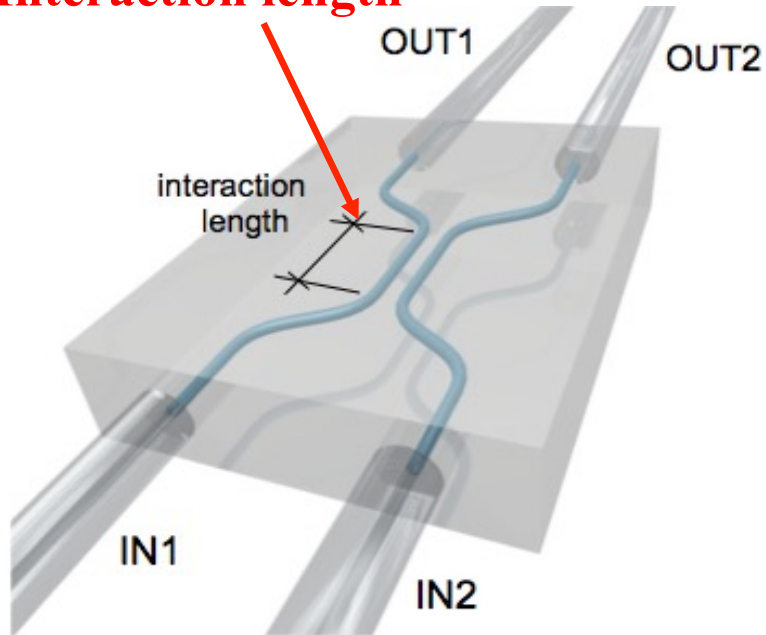


Polarizing beam splitter

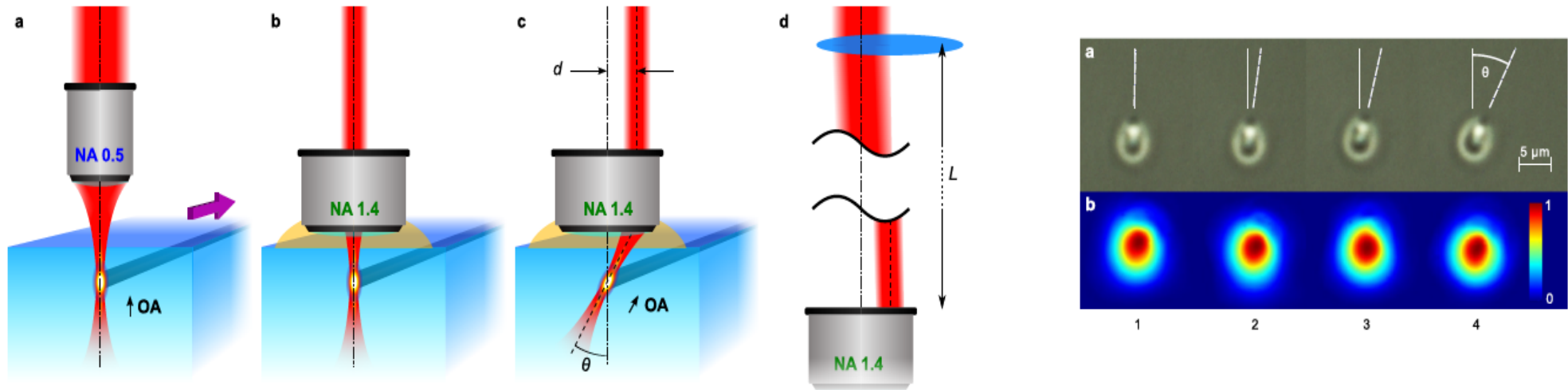
Different periodicities for H and V polarization deriving from residual asymmetry of waveguides.

Useful to realize polarization dependent devices (PPDC)

Interaction length

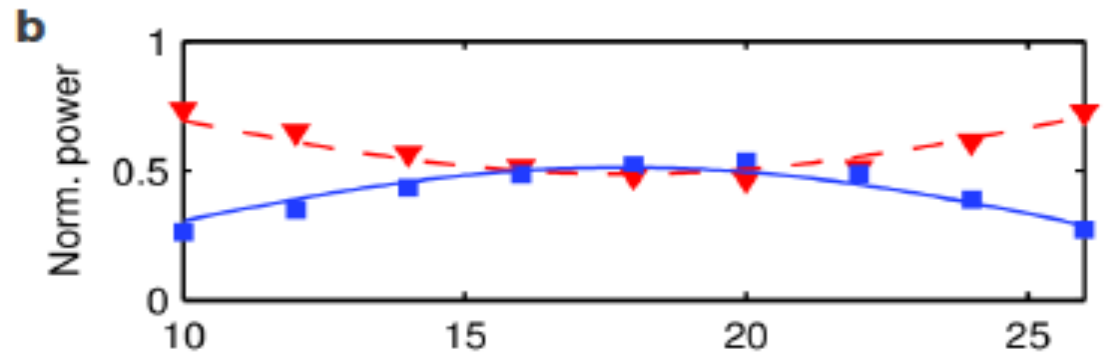


Integrated waveplates

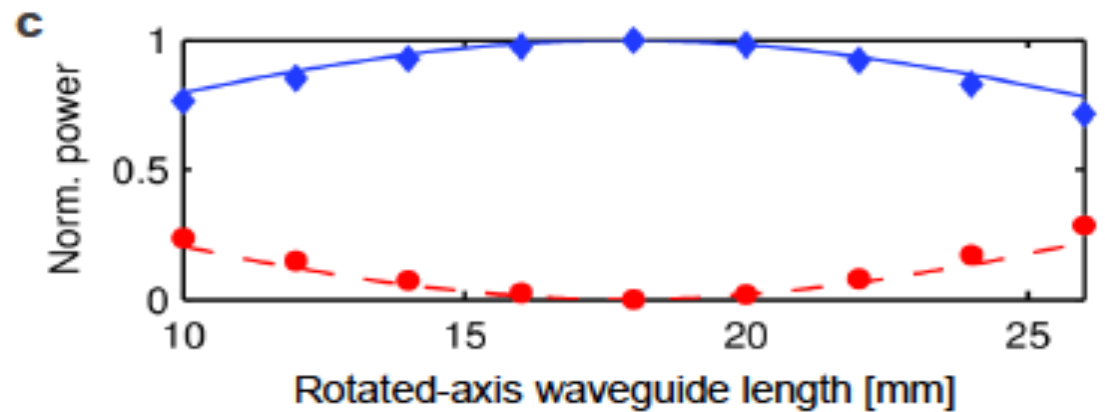


Fabrication

Input: H/V



Input: D/A



2-photon hyperentanglement

Entangling two photons in many degrees of freedom: alternative to distributing the qubits between many particles (multiphoton entanglement)

$$|\Psi_N\rangle = |Bell_1\rangle \otimes |Bell_2\rangle \dots \otimes |Bell_N\rangle$$

$N = \#$ of degrees of freedom

$|Bell_i\rangle :$

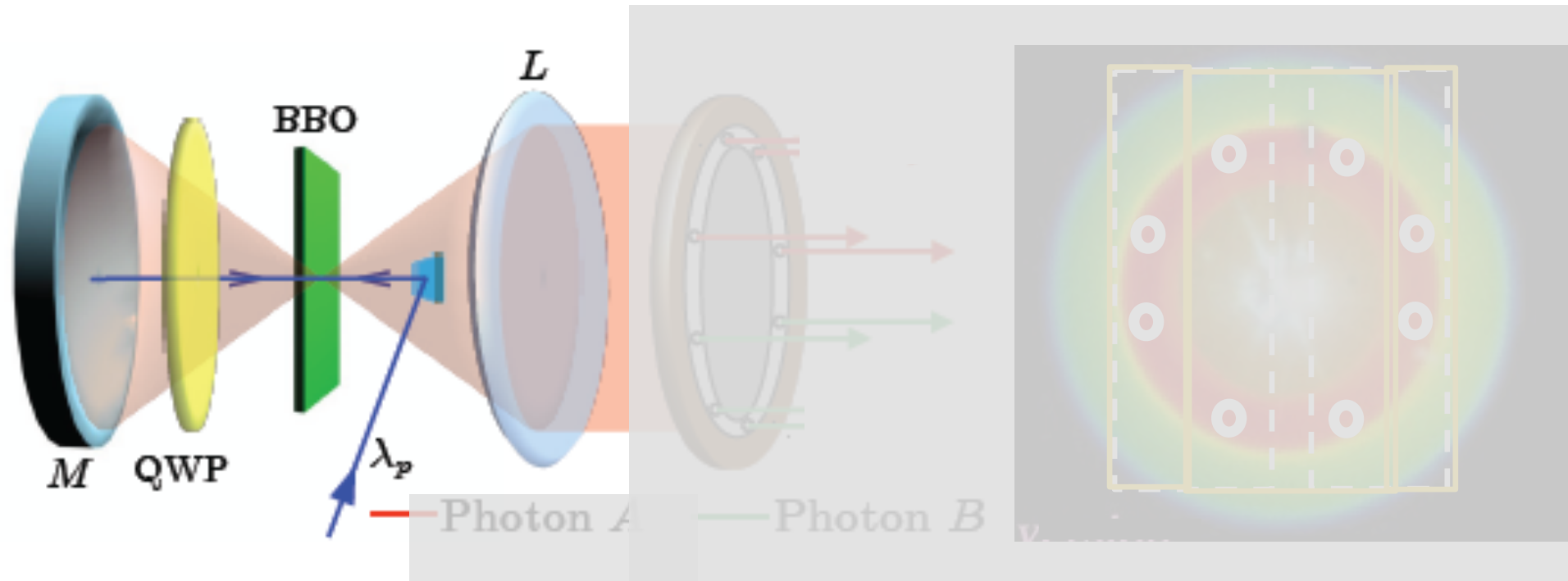
$$|\Phi^\pm\rangle = \frac{1}{\sqrt{2}} [|00\rangle \pm |11\rangle]$$

$$|\Psi^\pm\rangle = \frac{1}{\sqrt{2}} [|01\rangle \pm |10\rangle]$$

- Less decoherence with respect to n -photon entanglement.
- Detection efficiency, scaling as η^N , is constant when growing the size of the state.
- Exponential growing of resources but higher repetition in the state detection/generation rate.
- Used for advanced tasks of quantum information : Bell State Analysis, superdense coding, one-way quantum computation, advanced tests of quantum nonlocality, multipartite entangled states...

Path-polarization hyperentanglement

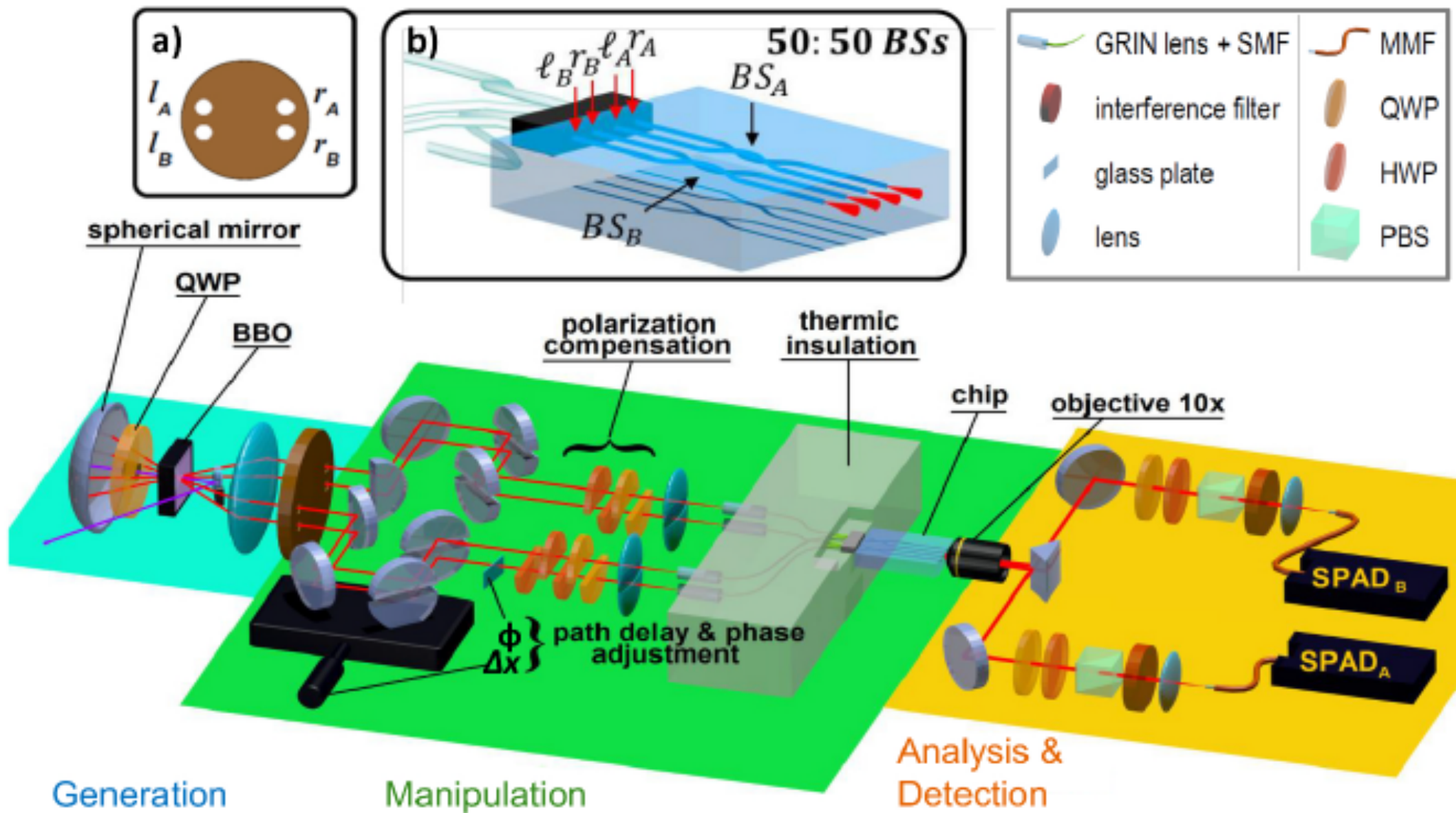
Use path and polarization of 2 photons to create 4/6 qubits)



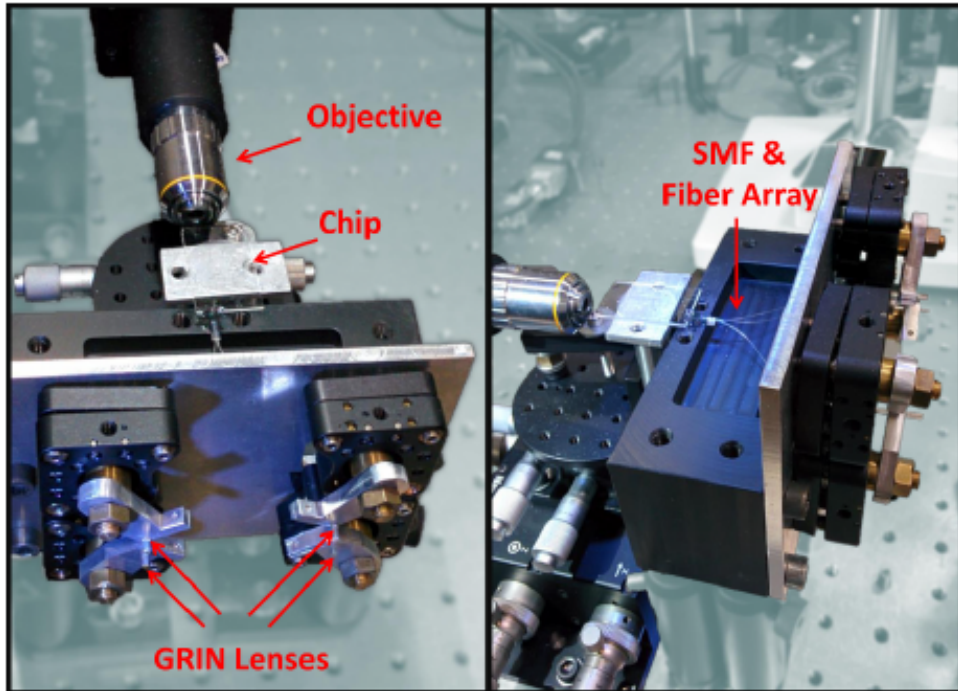
$$HE_6 = \frac{1}{\sqrt{2}} [|H_A H_B\rangle - |V_A V_B\rangle] \otimes \frac{1}{\sqrt{2}} [|l_A r_B\rangle + |r_A l_B\rangle] \otimes \frac{1}{\sqrt{2}} [|E_A E_B\rangle + |I_A I_B\rangle]$$

2 photons → 4 qubits

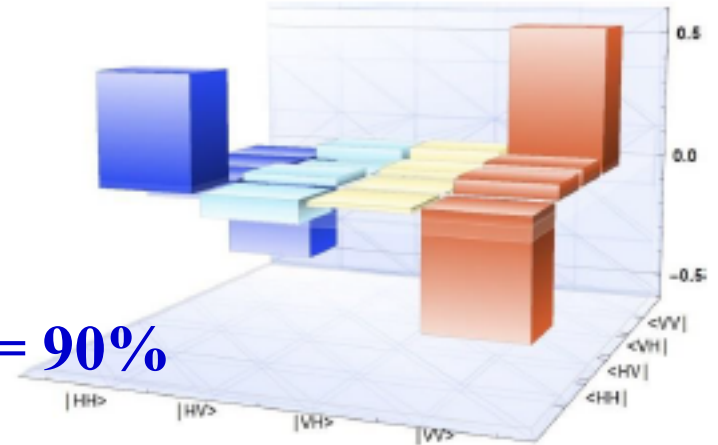
Hyperentanglement on a chip



Hyperentanglement on a chip

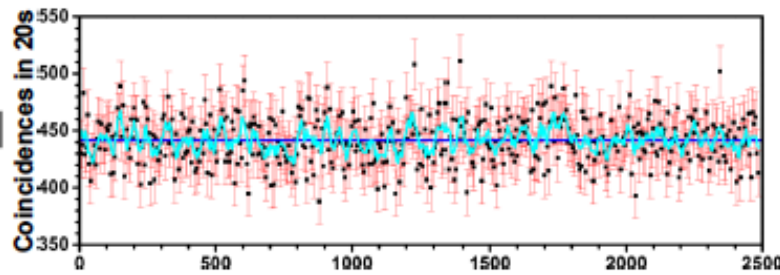
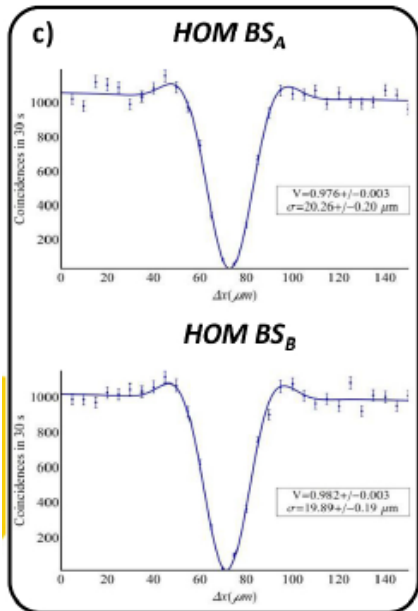


$F = 90\%$



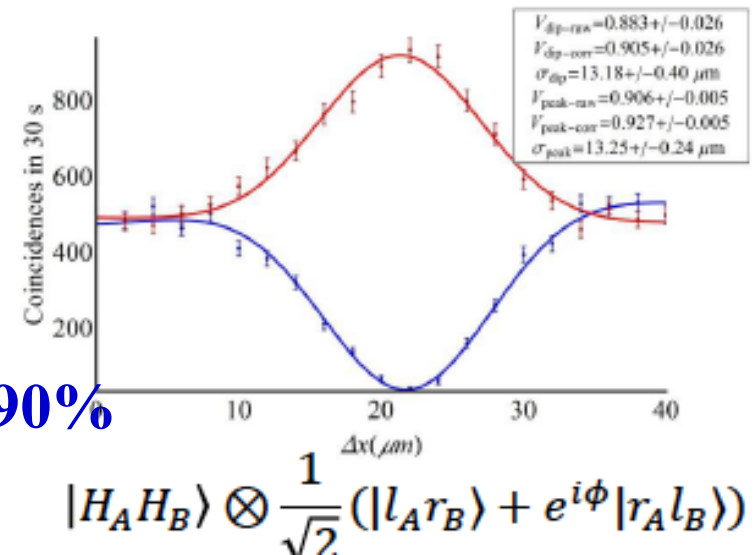
$$|l_A r_B\rangle \otimes \frac{1}{\sqrt{2}} (|H_A H_B\rangle - |V_A V_B\rangle)$$

Test separately polarization and path entanglement



Good phase stability

$V > 90\%$

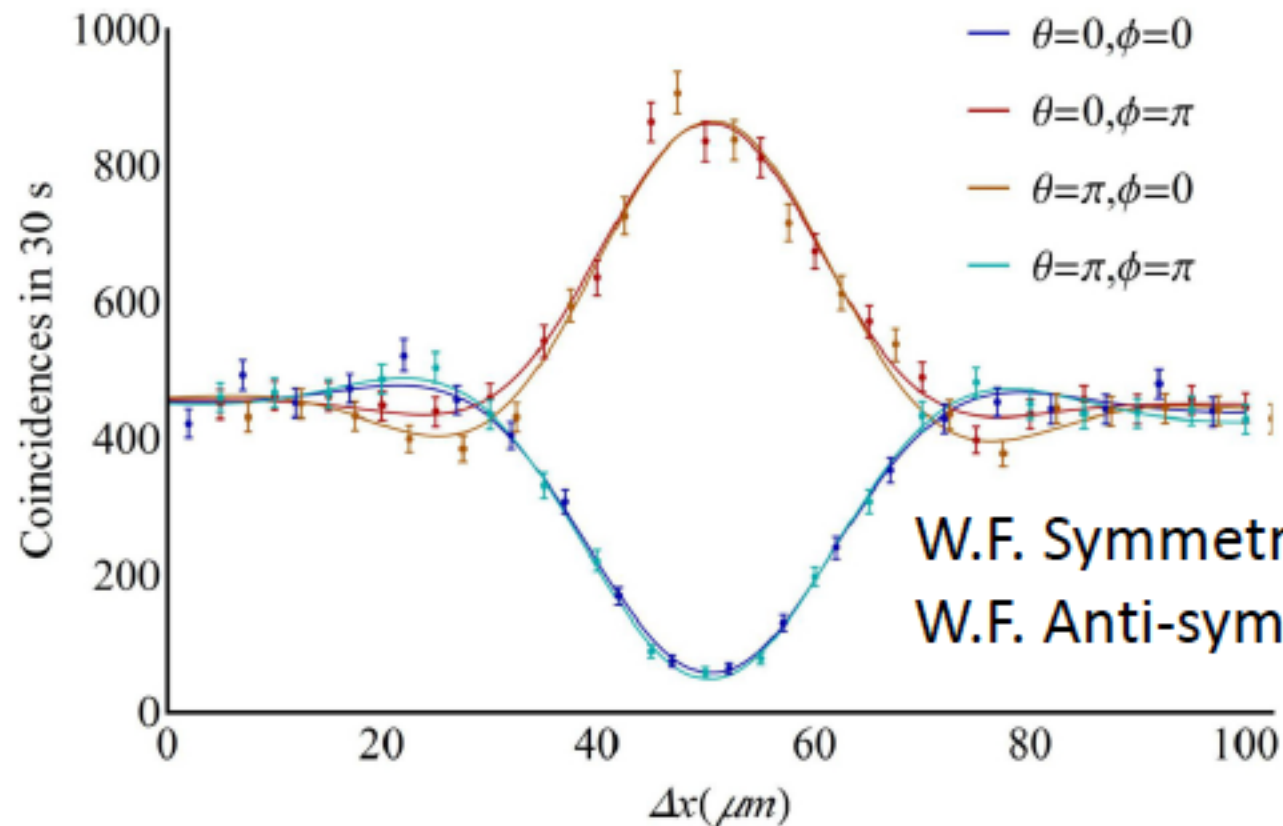


$$|H_A H_B\rangle \otimes \frac{1}{\sqrt{2}} (|l_A r_B\rangle + e^{i\phi} |r_A l_B\rangle)$$

Hyperentanglement measurements

Hyperentangled state interference:

$$|\Omega\rangle = \frac{1}{2} (|HH\rangle_{AB} + e^{i\phi} |VV\rangle_{AB}) \otimes (|rl\rangle_{AB} + e^{i\theta} |lr\rangle_{AB})$$

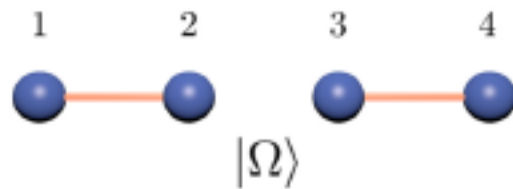


W.F. Symmetric -> **bosonic** behaviour

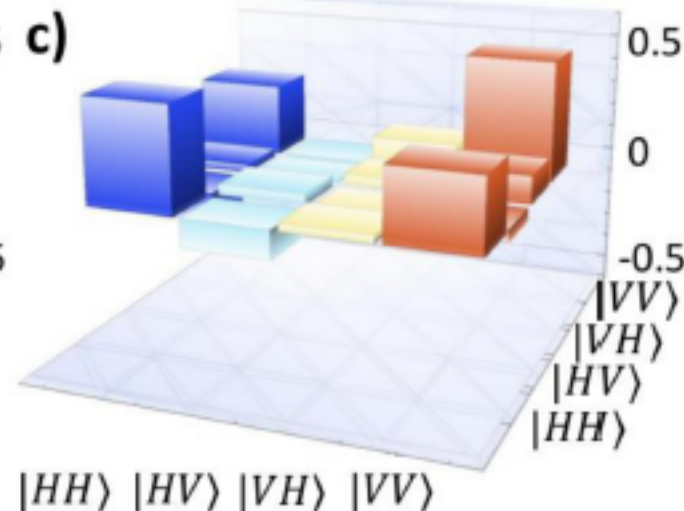
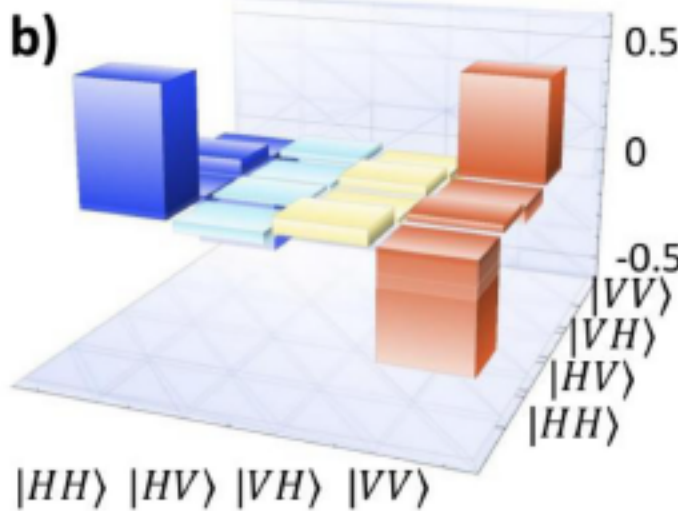
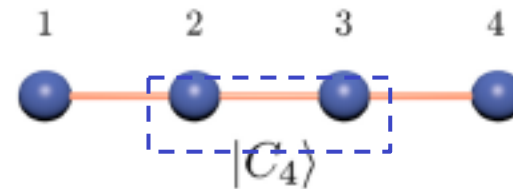
W.F. Anti-symmetric -> **fermionic** behaviour

Hyperentangled cluster states

a) Hyperentangled State



Cluster State



$$F_{\Phi^-} = 0.91 \pm 0.10,$$

$$C_{\Phi^-} = 0.88 \pm 0.08,$$

$$F_{\Phi^+} = 0.83 \pm 0.11,$$

$$C_{\Phi^+} = 0.91 \pm 0.08.$$

$$|C_4\rangle = \frac{1}{2}(|H_a r_A H_B l_B\rangle + |V_a r_A B_B l_B\rangle + |H_a l_A H_B r_B\rangle - |V_a l_A V_B r_B\rangle)$$

$|C_4\rangle$ realized by a HW on mode r_A acting as a Control Phase (CP).
90% fidelity measured for the states ψ_+ and ψ_- on pairs $l_A r_B$ and $r_A l_B$.

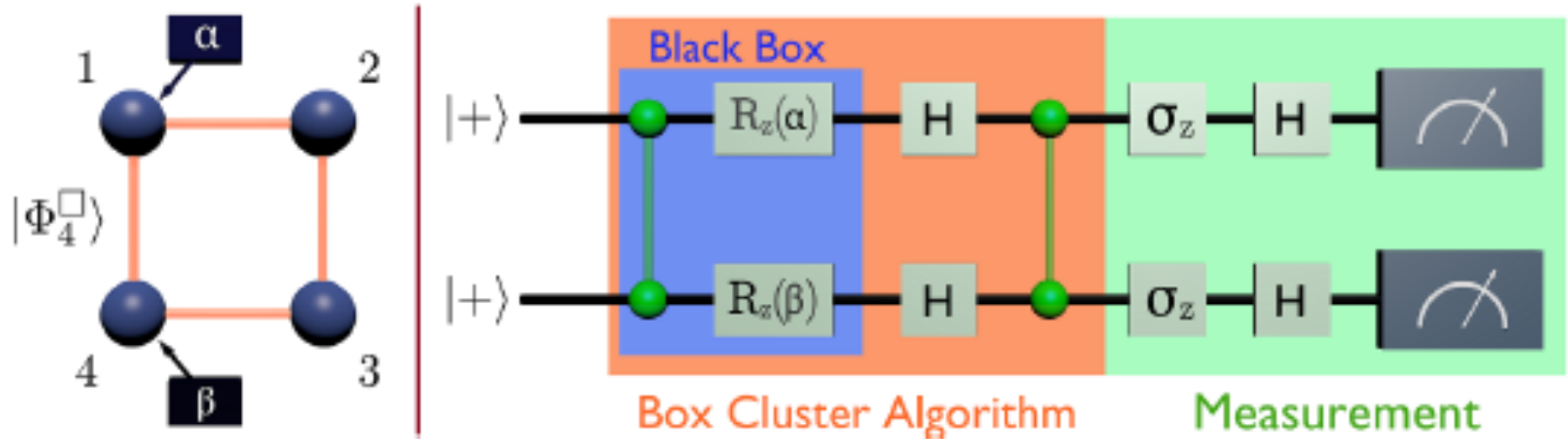
Cluster state measurements

$Z_A Z_B$	$+0.940 \pm 0.028$
$X_A X_B z_A$	-0.860 ± 0.030
$X_A X_B z_B$	$+0.860 \pm 0.030$
$z_A z_B$	-0.990 ± 0.007
$Z_A x_A x_B$	$+0.8092 \pm 0.036$
$Z_B x_A x_B$	$+0.8081 \pm 0.035$

Use multipartite entanglement witness:

$$W = 0.5(4 - Z_A Z_B - Z_A x_A x_B + X_A z_A X_B + z_A z_B - x_A Z_B x_B - X_A X_B z_B) = -0.634 \pm 0.036$$

Grover's search algorithm



Have 2^M elements encoded in M qubits and a black box (oracle) tagging one of them by changing its sign.

Goal: identify the tagged item by a repeated query to the black box.

Classical: $2^M/2$ calculations

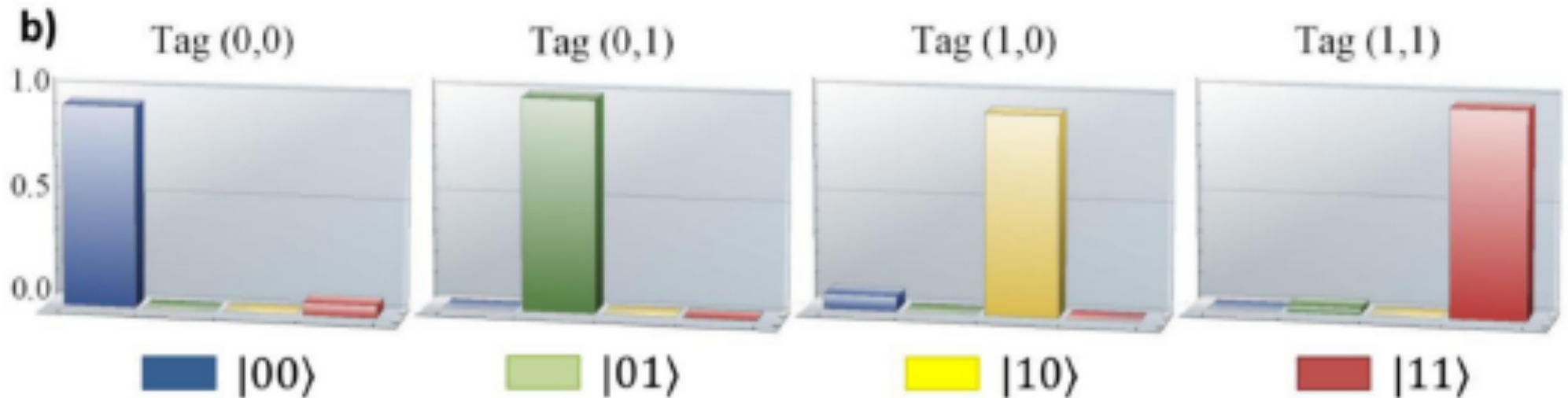
Grover: $\sqrt{2^M}$ operations

Four qubit box cluster states composed by the path (k) and polarization (π) qubits, labelled in the order $(1, 2, 3, 4) = (k_B, \pi_A, k_A, \pi_B)$.

Measure qubits 1 and 4 on the bases α and β .

Information processed and read on qubits 2 and 3.

Grover's algorithm: results



Outcome probability for different tagged items

Probabilistic algorithm (depends on postselection):

Average probability = 0.960 ± 0.007 (average rate: 17 Hz)

Deterministic algorithm (passive feed-forward)

Average probability = 0.964 ± 0.003 (average rate: 68 Hz)

Challenges, perspectives.....

- **Increase number of modes (limitation: bending losses of laser written waveguides)**
- **Active reconfiguration of Unitaries**
- **Novel 3D structures for more complex circuits**
- **Exponential growth of quantum complexity**
- **More complex circuits including PBSs, waveplates, active phase shifters enabling state manipulation and analysis on a chip**
- **Integrate hyperentanglement source inside the chip.**
- **High efficiency single photon detectors**
- **Higher number of photons**



Fabio Sciarrino



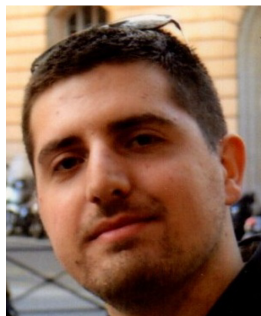
Nicolò Spagnolo
Postdoc



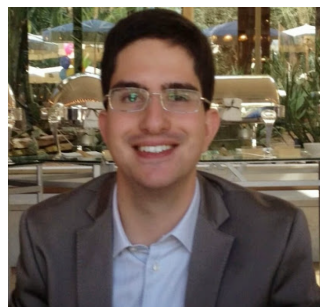
Marco Bentivegna
PhD student



Fulvio Flammini
PhD student



Niko Viggianiello
PhD student



Mario A. Ciampini
PhD student



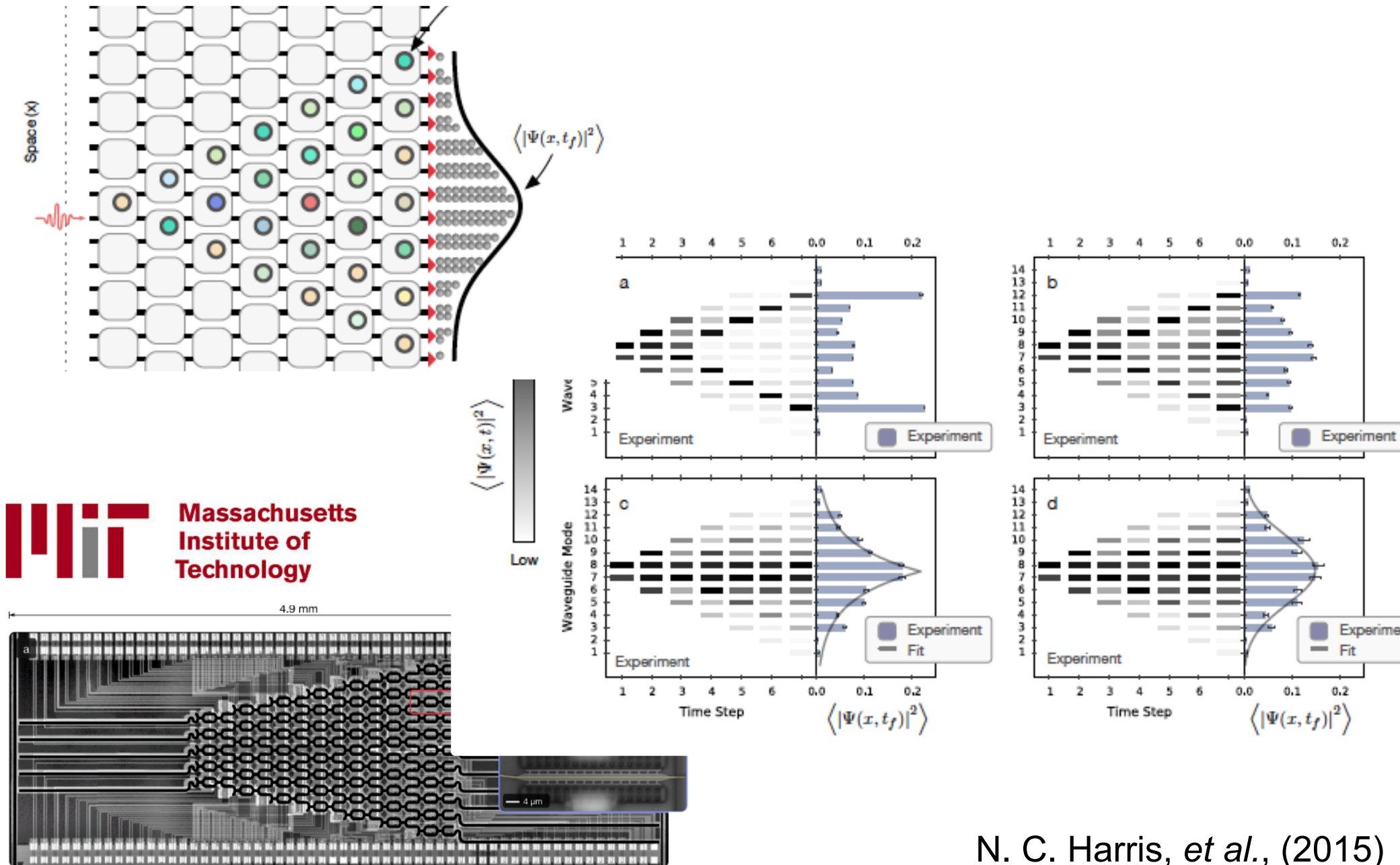
Adeline Orieux
Postdoc (now at
Telecom ParisTech)



Chiara Vitelli
Postdoc (now at
Authority
per l'Energia)



Bosonic transport simulation in a programmable processor



Universal linear optics

